

COUNTING PARTITIONS INSIDE A RECTANGLE*

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Abstract. We consider the number of partitions of n whose Young diagrams fit inside an $m \times \ell$ rectangle; equivalently, we study the coefficients of the q -binomial coefficient $\binom{m+\ell}{m}_q$. We obtain sharp asymptotics throughout the regime $\ell = \Theta(m)$ and $n = \Theta(m^2)$, while previously sharp asymptotics were derived by Takács [31] only in the regime where $|n - \ell m/2| = O(\sqrt{\ell m(\ell + m)})$ using a local central limit theorem. Our approach is to solve a related large deviation problem: we describe the tilted measure that produces configurations whose bounding rectangle has the given aspect ratio and is filled to the given proportion. Our results are sufficiently sharp to yield the first asymptotic estimates on the consecutive differences of these numbers when n is increased by one and m, ℓ remain the same, hence significantly refining Sylvester’s unimodality theorem and giving effective asymptotic estimates for related Kronecker and plethysm coefficients from representation theory.

Key words. partitions, q -binomial coefficients, Kronecker coefficients, central limit theorem.

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1. Introduction. A partition λ of n is a sequence of weakly decreasing nonnegative integers $\lambda = (\lambda_1 \geq \lambda_2 \geq \dots)$ whose sum $|\lambda| = \lambda_1 + \lambda_2 + \dots$ is equal to n . The study of integer partitions is a classic subject with applications ranging from number theory to representation theory and combinatorics, and integer partitions with various restrictions on properties, such as part sizes or number of parts, occupy the field of partition theory [2]. The generating functions of integer partitions play a role in number theory and the theory of modular forms. In representation theory, integer partitions index the conjugacy classes and irreducible representations of the symmetric group S_n ; they are also the signatures of the irreducible polynomial representation of GL_n and give a basis for the ring of symmetric functions. More recently, partitions have appeared in the study of interacting particle systems and other statistical mechanics models.

The number of partitions of n , typically denoted by $p(n)$ but here un conventionally¹ by N_n , was implicitly determined by Euler via the generating function

$$\sum_{n=0}^{\infty} N_n q^n = \prod_{i=1}^{\infty} \frac{1}{1 - q^i}.$$

There is no exact explicit formula for the numbers N_n . The asymptotic formula

$$(1.1) \quad N_n := \#\{\lambda \vdash n\} \sim \frac{1}{4n\sqrt{3}} \exp\left(\pi\sqrt{\frac{2n}{3}}\right),$$

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¹We use the notation N_n to distinguish scenarios of probability with those of enumeration, both of which occur in the present manuscript.

34 obtained by Hardy and Ramanujan [10], is considered to be the beginning of the use
 35 of complex variable methods for asymptotic enumeration of partitions (the so-called
 36 circle method).

37 Our goal is to obtain asymptotic formulas similar to (1.1) for the number of
 38 partitions λ of n whose Young diagram fits inside an $m \times \ell$ rectangle, denoted

$$39 \quad N_n(\ell, m) := \#\{\lambda \vdash n : \lambda_1 \leq \ell, \text{ length}(\lambda) \leq m\}.$$

40 These numbers are also the coefficients in the expansion of the q -binomial coefficient

$$41 \quad \binom{\ell + m}{m}_q = \frac{\prod_{i=1}^{\ell+m} (1 - q^i)}{\prod_{i=1}^{\ell} (1 - q^i) \prod_{i=1}^m (1 - q^i)} = \sum_{n=0}^{\ell m} N_n(\ell, m) q^n.$$

42 The q -binomial coefficients are themselves central to enumerative and algebraic
 43 combinatorics. They are the generating functions for lattice paths restricted to rec-
 44 tangles and taking only north and east steps under the area statistic, given by the
 45 parameter n . They are also the number of ℓ -dimensional subspaces of $\mathbb{F}_q^{\ell+m}$ and
 46 appear in many other generating functions as the q -analogue generalization of the
 47 ubiquitous binomial coefficients. Notably, the numbers $N_n(\ell, m)$ form a symmetric
 48 unimodal sequence

$$49 \quad 1 = N_0(\ell, m) \leq N_1(\ell, m) \leq \dots \leq N_{\lfloor m\ell/2 \rfloor}(\ell, m) \geq \dots \geq N_{m\ell}(\ell, m) = 1,$$

50 a fact conjectured by Cayley in 1856 and proven by Sylvester in 1878 via the repre-
 51 sentation theory of sl_2 [26]. One hundred forty years later, no previous asymptotic
 52 methods have been able to prove this unimodality.

53 **Asymptotics of $N_n(\ell, m)$.** Our first result is an asymptotic formula for $N_n(\ell, m)$
 54 in the regime $\ell/m \rightarrow A$ and $n/m^2 \rightarrow B$ for any fixed $A > B > 0$. This is the regime
 55 in which a limit shape of the partitions exists: $\ell/m \rightarrow A$ implies the aspect ratio has
 56 a limit and $n/m^2 \rightarrow B \in (0, A)$ implies the portion of the $m \times \ell$ rectangle that is filled
 57 approaches a value that is neither zero nor one. By ‘‘asymptotic formula’’ we mean
 58 a formula giving $N_n(\ell, m)$ up to a factor of $1 + o(1)$; such asymptotic equivalence is
 59 denoted with the symbol \sim . By replacing a partition with its complements in an $\ell \times m$
 60 rectangle, one sees that $N_n(\ell, m) = N_{m\ell-n}(\ell, m)$ and it thus suffices to consider only
 61 the case $A \geq 2B > 0$.

62 To state our results, given $A \geq 2B > 0$ we define three quantities c, d and Δ .
 63 The quantities c and d are the unique positive real solutions (see Lemma 9) to the
 64 simultaneous equations

$$65 \quad (1.2) \quad A = \int_0^1 \frac{1}{1 - e^{-c-dt}} dt - 1 = \frac{1}{d} \log \left(\frac{e^{c+d} - 1}{e^c - 1} \right) - 1,$$

$$66 \quad (1.3) \quad B = \int_0^1 \frac{t}{1 - e^{-c-dt}} dt - \frac{1}{2} = \frac{d \log(1 - e^{-c-d}) + \text{dilog}(1 - e^{-c}) - \text{dilog}(1 - e^{-c-d})}{d^2},$$

67 where we recall the dilogarithm function

$$69 \quad \text{dilog}(x) = \int_1^x \frac{\log t}{1 - t} dt = \sum_{k=1}^{\infty} \frac{(1 - x)^k}{k^2}$$

70 for $|x - 1| < 1$. The quantity Δ , which will be seen to be strictly positive, is defined
 71 by

$$72 \quad (1.4) \quad \Delta = \frac{2Be^c(e^d - 1) + 2A(e^c - 1) - 1}{d^2(e^{d+c} - 1)(e^c - 1)} - \frac{A^2}{d^2}.$$

73 THEOREM 1. Given m, ℓ and n , let $A := \ell/m$ and $B := n/m^2$ and define c, d and
 74 Δ as above. Let K be any compact subset of $\{(x, y) : x \geq 2y > 0\}$. As $m \rightarrow \infty$ with
 75 ℓ and n varying so that (A, B) remains in K ,

$$76 \quad (1.5) \quad N_n(\ell, m) \sim \frac{e^{m[cA+2dB-\log(1-e^{-c-d})]}}{2\pi m^2 \sqrt{\Delta(1-e^{-c})(1-e^{-c-d})}},$$

77 where c and d vary in a Lipschitz manner with $(A, B) \in K$.

REMARK. In the special case $B = A/2$, the parameters take on the elementary values

$$d = 0, \quad c = \log\left(\frac{A+1}{A}\right), \quad \text{and} \quad \Delta = \frac{A^2(A+1)^2}{12}.$$

In this case we understand the exponent and leading constant to be their limits as $d \rightarrow 0$, giving

$$N_{Am^2/2}(Am, m) \sim \frac{\sqrt{3}}{A\pi m^2} \left[\frac{(A+1)^{A+1}}{A^A} \right]^m.$$

78 The special case when $A \rightarrow \infty$, so that $N_n(\ell, m) = N_n(m)$ and the restriction
 79 on partition sizes is removed, corresponds to taking $c = 0$ and having d be a solu-
 80 tion to an explicit equation given in Lemma 9. In this case the result matches the
 81 one obtained first by Szekeres [29] using complex analysis, then by Canfield [5] using
 82 a recursion, and most recently by Romik [21] using probabilistic methods based on
 83 Fristedt's ensemble [9]. These works and others are further explained in Section 2.

84 **Unimodality.** Our second result gives an asymptotic estimate of the consecu-
 85 tive differences of N_n . In fact our motivation for deriving more accurate asymptotics
 86 for $N_n(\ell, m)$ was to be able to analyze the sequence $\{N_{n+1}(\ell, m) - N_n(\ell, m) : n \geq$
 87 $1\}$. Sylvester's proof of unimodality of $N_n(\ell, m)$ in n [26], and most subsequent
 88 proofs [23, 24, 19], are algebraic, viewing $N_n(\ell, m)$ as dimensions of certain vector
 89 spaces, or their differences as multiplicities of representations. While there are also
 90 purely combinatorial proofs of unimodality, notably O'Hara's [14] and the more ab-
 91 stract one in [18], they do not give the desired symmetric chain decomposition of the
 92 subposet of the partition lattice. These methods do not give ways of estimating the
 93 asymptotic size of the coefficients or their difference. It is now known that $N_n(\ell, m)$
 94 is strictly unimodal [15], and the following lower bound on the consecutive difference
 95 was obtained in [16, Theorem 1.2] using a connection between integer partitions and
 96 Kronecker coefficients:

$$97 \quad (1.6) \quad N_n(\ell, m) - N_{n-1}(\ell, m) \geq 0.004 \frac{2^{\sqrt{s}}}{s^{9/4}},$$

98 where $n \leq \ell m/2$ and $s = \min\{2n, \ell^2, m^2\}$. In particular, when $\ell = m$ we have $s = 2n$.

99 Any sharp asymptotics of the difference appears to be out of reach of the algebraic
 100 methods in this previous series of papers. Refining Theorem 1, we are able to obtain
 101 the following estimate.

THEOREM 2. Given m, ℓ and n , let $A := \ell/m$ and $B := n/m^2$ and define d as above. Suppose m, ℓ and n go to infinity so that (A, B) remains in a compact subset K of $\{(x, y) : x \geq 2y > 0\}$ and

$$m^{-1} |n - \ell m/2| \rightarrow \infty.$$

Then

$$N_{n+1}(\ell, m) - N_n(\ell, m) \sim \frac{d}{m} N_n(\ell, m).$$

102 REMARK. The condition $m^{-1} |n - lm/2| \rightarrow \infty$ is equivalent to $m |A - B/2| \rightarrow \infty$
 103 and is satisfied if and only if d , which depends on m , is not $O(m^{-1})$. It is automati-
 104 cally satisfied whenever the compact set K is a subset of $\{(x, y) : x > 2y > 0\}$.

105 **Corollary: Asymptotics of Kronecker coefficients.** Recent developments
 106 in the representation theory of the symmetric and general linear groups, motivated
 107 by applications to Computational Complexity theory, have realized the consecutive
 108 differences $N_{n+1}(\ell, m) - N_n(\ell, m)$ as specific Kronecker coefficients for the tensor
 109 product of irreducible $S_{\ell m}$ representations (see, for instance, [15] which is also one of
 110 the unimodality proofs). The Kronecker coefficient $g(\lambda, \mu, \nu) = \dim \text{Hom}(\mathbb{S}_\lambda, \mathbb{S}_\mu \otimes \mathbb{S}_\nu)$
 111 is the multiplicity of the irreducible $S_{|\lambda|}$ Specht module \mathbb{S}_λ in the tensor product of
 112 two other irreducible representations. It is a notoriously hard problem to determine
 113 the values of these coefficients, and their combinatorial interpretation has been an
 114 outstanding open problem in Algebraic Combinatorics since their definition by Mur-
 115 naghan in 1938 (see Stanley [25]). In general, determining even whether Kronecker
 116 coefficients are nonzero is an NP-hard problem and it is not known whether computing
 117 them lies in NP. See [11] and the literature therein for some recent developments on
 118 the relevance of Kronecker coefficients in distinguishing complexity classes on the way
 119 towards $P \neq NP$. Being able to estimate particular values of Kronecker coefficients is
 120 crucial to the Geometric Complexity Theory approach towards these problems.

121 Because it is known [15] that the consecutive difference $N_n(m, \ell) - N_{n-1}(m, \ell)$
 122 equals the Kronecker coefficient $g((m\ell - n, n), (m^\ell), (m^\ell))$, Theorem 2 gives the first
 123 tight asymptotic estimate on this family of Kronecker coefficients.

124 COROLLARY 3. The Kronecker coefficient of $S_{m\ell}$ for the (rectangle, rectangle,
 125 two-row) case is asymptotically given by

$$\begin{aligned} 126 \quad g((m^\ell), (m^\ell), (m\ell - n - 1, n + 1)) &= N_{n+1}(\ell, m) - N_n(\ell, m) \\ 127 \quad &\sim \frac{d}{m} N_n(\ell, m) \\ 128 \quad &\sim \frac{de^{m[cA + 2dB - \log(1 - e^{-c-d})]}}{2\pi m^3 \sqrt{\Delta(1 - e^{-c})(1 - e^{-c-d})}}, \\ 129 \end{aligned}$$

130 with constants and ranges as in Theorems 1 and 2.

131 2. Review of previous results and description of methods.

132 **2.1. Combinatorial Enumeration.** Work on this problem has developed in
 133 two streams. First, there have been combinatorial results aimed at asymptotic enu-
 134 meration in various regimes. After Hardy and Ramanujan obtained an asymptotic
 135 formula for N_n in [10], enumerative work focused on $N_n(m)$, the number of partitions
 136 with part sizes bounded by m , or equivalently, partitions of n that fit in an $m \times \infty$
 137 strip of growing height. In 1941, Erdős and Lehner [8] showed that $N_n(m) \sim \frac{n^{m-1}}{m!(m-1)!}$
 138 for $m = o(n^{1/3})$. This was generalized by Szekeres and others, culminating in asymp-
 139 totics of $N_n(m)$ for all m in 1953 [29]. Szekeres simplified his arguments a number
 140 of times, ultimately giving asymptotics using only a saddle-point analysis, without
 141 needing results on modular functions; his argument has been referred to as the Szek-
 142 eres circle method. Canfield [5] gave an elementary proof (with no complex analysis)
 143 of asymptotics for $N_n(m)$ using a recursive formula satisfied by these numbers.

The combinatorial stream contains a few results on the asymptotics of $N_n(m, \ell)$ but only in the regime where m and ℓ are greater than \sqrt{n} by at least a factor of $\log n$. This is a natural regime to study because the typical values of the maximum part (equivalently the number of parts) of a partition of size n was shown by Erdős and Lehner [8] to be of order $\sqrt{n \log n}$. Szekeres [30, Theorem 1] used saddle-point techniques to express $N_n(\ell, m)$ in terms of N_n , $\lambda := \frac{\pi \ell}{\sqrt{6n}}$ and $\mu := \frac{\pi m}{\sqrt{6n}}$. If, in fact,

$$\frac{\sqrt{6n}}{\pi} \left(\frac{1}{4} + \varepsilon \right) \log n < \ell, m < \frac{\sqrt{6n} \log n}{\pi}$$

for some $\varepsilon > 0$, then the distributions defined by ℓ and m are independent and equal, and Szekeres' formula simplifies to

$$N_n(\ell, m) \sim N_n \exp \left[-(\lambda + \mu) - \sqrt{\frac{6n}{\pi}} (e^{-\lambda} + e^{-\mu}) \right].$$

144 The Szekeres circle method was recently revisited by Richmond [20]. In [12] the au-
 145 thors, independently and concurrently with our paper, used the generating function
 146 for q -binomial coefficients and a saddle point analysis to derive the asymptotics for
 147 $N_n(m, \ell)$ in the cases when $m, \ell \geq 4\sqrt{n}$, corresponding to $B \leq \min\{1, A^2\}/16$ in our
 148 notation. Those authors express their result using the root of a hypergeometric iden-
 149 tity similar to (1.3), however their methods give weaker error bounds and consequently
 150 cannot answer questions of unimodality.

151 **2.2. Probabilistic limit theorems.** The second strand of work on this problem
 152 has been probabilistic. The goal in this strand has been to determine properties of
 153 a random partition or Young diagram, picked from a suitable probability measure.
 154 This approach goes back at least to Mann and Whitney [13], who showed that the
 155 size of a uniform random partition contained in an $\ell \times m$ rectangle satisfies a normal
 156 distribution. Fristedt [9] defined a distribution on partitions of all sizes, weighted
 157 with respect to a parameter $q < 1$. The key property of the measure employed is
 158 that it makes the number $X_k(\lambda)$ of parts of size k in the partition λ drawn under
 159 this distribution independent as k varies; the distributions of the X_k are reduced
 160 geometrics with respective parameter $1 - q^k$, so that their mean is $q^k / (1 - q^k)$. Fristedt
 161 is chiefly concerned with the limiting behavior of kX_k for $k = o(\sqrt{n})$, which rescales,
 162 on division by \sqrt{n} , to an exponential distribution. A line of work beginning with
 163 Sinaĭ [22] uses similar methods to study convex polygons with various restrictions. In
 164 particular, Sinaĭ defines a distribution on convex polygons which is uniform on walks
 165 with fixed endpoints, then tunes parameters of the distribution so that a local limit
 166 theorem holds. More recent work of Bureaux [4] continues this approach to study
 167 partitions of two-dimensional integer vectors.

168 Much of the work following Fristedt's is concerned with a description of the lim-
 169 iting shape of the random partition, and fluctuations around that shape. The limit
 170 shape of an unrestricted partition was posed as a problem by Vershik and first an-
 171 swered in [27, 28]. In 2001, Vershik and Yakubovich [32] describe the limit shape for
 172 singly restricted partitions: those with $m \leq c\sqrt{n}$. They obtain both main (strong
 173 law) results and fluctuation (CLT) results. It is in this paper that the probability
 174 measures \mathbb{P}_m used in our analysis below first arose, although we were unaware of this
 175 when we first derived them from large deviation principles. The limit shape for doubly
 176 restricted partitions in the regime $m, \ell = \Theta(\sqrt{n})$ was first described by Petrov [17].
 177 It is identified there with a portion of the curve $e^{-x} + e^{-y} = 1$, which represents the

178 limit shape of unrestricted partitions. More recently, Beltoft et al. [3] obtained fluctuation results in the doubly restricted regime. The limiting fluctuation process is an
 179 Ornstein-Uhlenbeck bridge, generalizing the two-sided stationary Ornstein-Uhlenbeck
 180 process that gives the limiting fluctuations in the unrestricted case [32].
 181

182 **2.3. Enumeration via probability.** Strangely, we know of only one paper
 183 combining these two streams. Takács [31] observed the following consequence of the
 184 work of Fristedt and others. Begin a discrete walk at $(\ell, 0)$ and randomly choose
 185 steps in the $(0, -1)$ or $(-1, 0)$ directions by making independent fair coin flips. If this
 186 walk goes from $(\ell, 0)$ to $(0, -m)$ it takes precisely $m + \ell$ steps and encloses a Young
 187 diagram fitting in an $m \times \ell$ rectangle: see Figure 1. Let $G(m, \ell)$ denote the event that
 188 a walk of length $m + \ell$ ends at $(0, -m)$ and let $H(m, n)$ denote the event that the
 189 resulting Young diagram has area n . Under the IID fair coin flip probability measure
 190 on paths, all paths of length $m + \ell$ have the same probability $2^{-(m+\ell)}$. Therefore,
 191 $\mathbb{P}[G(m, \ell) \cap H(m, n)] = 2^{-(m+\ell)} N_n(\ell, m)$ and the problem of counting $N_n(\ell, m)$ is
 192 reduced to determining the probability $\mathbb{P}[G(m, \ell) \cap H(m, n)]$.

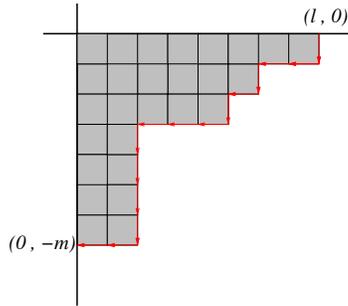


Fig. 1: The red arrows are the steps in a South and West directed simple random walk

193 Takács observed that this probability is computable by a two-dimensional local
 194 central limit theorem, ultimately obtaining bounds on the relative error that are of
 195 order $(m + \ell)^{-3}$. These error bounds are meaningful when n differs from $m\ell/2$ by
 196 up to a few multiples of $\log(m + \ell)$ standard deviations: if $\ell = \theta(m)$ this means that
 197 $|B - A/2|m^2 = \Theta(m^{3/2} \log m)$. When $|B - A/2| \gg m^{-1/2} \log m$ the error is much
 198 bigger than the main term of the Gaussian estimate provided by the LCLT and one
 199 cannot recover meaningful information about $N_n(\ell, m)$. This is where Takács left off
 200 and the present manuscript picks up.

201 **2.4. Description of our methods.** We use a local large deviation computation
 202 in place of a local central limit theorem: this is possible because the restriction to
 203 an $m \times \ell$ rectangle is a linear constraint. Indeed, consider now a partition $\lambda =$
 204 $(\lambda_1, \dots, \lambda_m)$ with at most m parts (so some λ_j may be zero) and define $\lambda_0 := \ell$ and
 205 $\lambda_{m+1} := 0$. It is convenient to encode a partition with respect to its *gaps* $x_j :=$
 206 $\lambda_j - \lambda_{j+1}$, so the condition that λ be a partition of n of size at most ℓ is equivalent
 207 to $x_j \geq 0$ and

$$208 \quad (2.1) \quad \sum_{j=0}^m x_j = \ell, \quad \sum_{j=0}^m jx_j = n.$$

209 Figure 2 gives a pictorial proof.

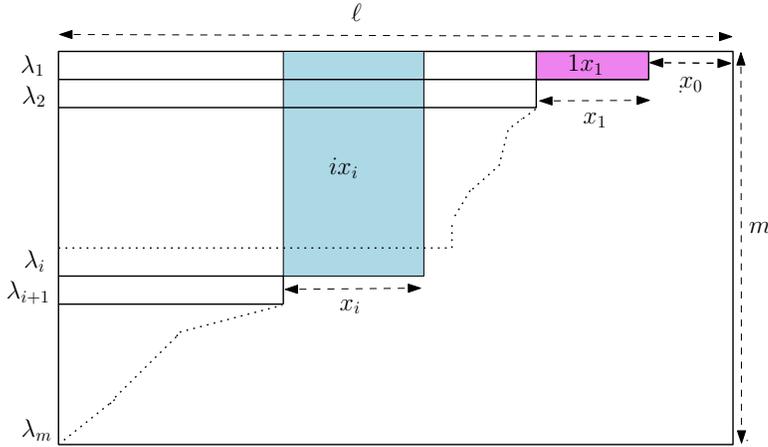


Fig. 2: The total area n of a partition is composed of rectangles of area jx_j

210 Solving the large deviation problem produces a “tilted measure” in which the
 211 gaps X_j are no longer IID reduced geometrics with parameter $1/2$ but are instead
 212 given by independent reduced geometric variables whose parameters $q_j = 1 - p_j$ vary
 213 in a log-linear manner. Log-linearity is dictated by the variational large deviation
 214 problem and leads to the same simplification as before. Not all partitions have the
 215 same probability under the tilted measure, but all those resulting in a given value of
 216 ℓ and n do have the same probability. Lastly, one must choose the particular linear
 217 function $\log q_j = -c - d(j/m)$ to ensure that λ being a partition of n with parts of size
 218 at most ℓ will again be in the central part of the tilted measure, so that asymptotics
 219 can be read off from a local CLT for the tilted measure.

220 The tilted measures \mathbb{P}_m that we employ are denoted $\mu_{x,y}$ in [32] and referred to
 221 there as the grand ensemble of partitions. That paper, however, was not concerned
 222 with enumeration, only with limit shape results. For this reason they do not state or
 223 prove enumeration results. In fact [17] is able to prove the shape result by estimating
 224 exponential rates only, showing rather elegantly that an ε error in the rescaled shape
 225 leads to an exponential decrease in the number of partitions. The present manuscript
 226 combines the idea of the grand ensemble with some precise central limit estimates and
 227 some algebra inverting the relation between the log-linear parameters and the param-
 228 eters A and B defining the respective limits of ℓ/m and n/m^2 to give estimates on
 229 $N_n(\ell, m)$ precise enough to also yield asymptotic estimates on $N_{n+1}(\ell, m) - N_n(\ell, m)$.

230 The first step of carrying this out necessarily recovers the leading exponential
 231 behavior for $N_n(\ell, m)$, which is implicit in [32] and [17] though Petrov only states it as
 232 an upper bound. Interestingly, Takács did not seem to be aware of the ease with which
 233 the exponential rate may be obtained. His result states a Gaussian estimate and an
 234 error term. As noted above, it is nontrivial only when the $(m+\ell)^{-3}$ relative error term
 235 does not swamp the main terms, which occurs when n is close to $\ell m/2$ (see also [1]).
 236 Figure 3 shows Takács’ predicted exponential growth rate on a family of examples
 237 compared to the actual exponential growth rate that follows from Theorem 1.

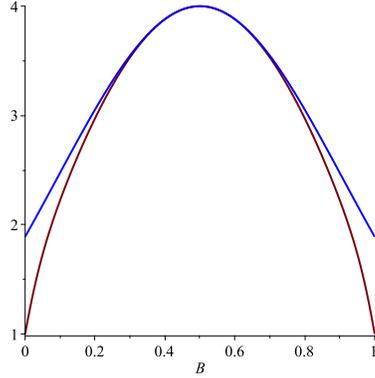


Fig. 3: Exponential growth of $N_{Bm^2}(m, m)$ predicted by Takács' formula (blue, above) compared to the actual exponential growth given by Theorem 1 (red, below).

238 **3. A discretized analogue to Theorem 1.** We now implement this program
 239 to derive asymptotics. With c_m and d_m to be specified later, let $q_j := e^{-c_m - jd_m/m}$,
 240 let $p_j := 1 - q_j$ and let

$$241 \quad L_m := \sum_{j=0}^m \log p_j.$$

242 Let \mathbb{P}_m be a probability law making the random variables $\{X_j : 0 \leq j \leq m\}$ indepen-
 243 dent reduced geometrics with respective parameters p_j . Define random variables S_m
 244 and T_m by

$$245 \quad (3.1) \quad S_m := \sum_{i=0}^m X_i; \quad T_m := \sum_{i=1}^m iX_i,$$

246 corresponding to the unique partition λ satisfying $X_j = \lambda_j - \lambda_{j+1}$. We first prove a
 247 result similar to Theorem 1, except that the parameters c and d that solve integral
 248 Equations (1.2) and (1.3) are replaced by c_m and d_m satisfying the discrete summation
 249 Equations (3.2) and (3.3) below. These equations say that $\mathbb{E}S_m = \ell$ and $\mathbb{E}T_m = m$.
 250 Writing this out, using $\mathbb{E}X_j = 1/p_j - 1 = 1/(1 - e^{-c_m - d_m j/m}) - 1$, gives

$$251 \quad (3.2) \quad \ell = \sum_{j=0}^m \frac{1}{1 - e^{-c_m - d_m j/m}} - (m+1)$$

$$252 \quad (3.3) \quad n = m \sum_{j=0}^m \frac{j/m}{1 - e^{-c_m - d_m j/m}} - \frac{m(m+1)}{2}.$$

254 Let M_m denote the covariance matrix for (S_m, T_m) . The entries may be computed

255 from the basic identity $\text{Var}(X_j) = q_j/p_j^2$, resulting in

$$256 \quad (3.4) \quad \text{Var}(S_m) = \sum_{j=0}^m \frac{e^{-c_m - d_m j/m}}{(1 - e^{-c_m - d_m j/m})^2}$$

$$257 \quad (3.5) \quad \text{Cov}(S_m, T_m) = \sum_{j=0}^m j \frac{e^{-c_m - d_m j/m}}{(1 - e^{-c_m - d_m j/m})^2}$$

$$258 \quad (3.6) \quad \text{Var}(T_m) = \sum_{j=0}^m j^2 \frac{e^{-c_m - d_m j/m}}{(1 - e^{-c_m - d_m j/m})^2}.$$

260 **THEOREM 4** (discretized analogue). *Let c_m and d_m satisfy (3.2) – (3.3). Define*
 261 *α_m, β_m and γ_m to be the normalized entries of the covariance matrix*

$$262 \quad \alpha_m := m^{-1} \text{Var}(S_m); \quad \beta_m := m^{-2} \text{Cov}(S_m, T_m); \quad \gamma_m := m^{-3} \text{Var}(T_m),$$

263 *which are $O(1)$ as $m \rightarrow \infty$. Again, let $A := \ell/m$ and $B := n/m^2$ and $\Delta_m :=$
 264 $\alpha_m \gamma_m - \beta_m^2$. Then as $m \rightarrow \infty$ with ℓ and n varying so that (A, B) remains in a
 265 compact subset of $\{(x, y) : x \geq 2y > 0\}$,*

$$266 \quad (3.7) \quad N_n(\ell, m) \sim \frac{1}{2\pi m^2 \sqrt{\Delta_m}} \exp \left\{ m \left(-\frac{L_m}{m} + c_m A + d_m B \right) \right\}.$$

267 *Proof.* The atomic probabilities $\mathbb{P}_m(\mathbf{X} = \mathbf{x})$ depend only on S_m and T_m as

$$268 \quad \log \mathbb{P}_m(\mathbf{X} = \mathbf{x}) = \sum_{j=0}^m (\log p_j + x_j \log q_j)$$

$$269 \quad = L_m - \sum_{j=0}^m \left(c_m + j \frac{d_m}{m} \right) x_j$$

$$270 \quad = L_m - c_m \left(\sum_{j=0}^m x_j \right) - \frac{d_m}{m} \left(\sum_{j=0}^m j x_j \right).$$

271
 272 In particular, for any \mathbf{x} satisfying (2.1),

$$273 \quad (3.8) \quad \log \mathbb{P}(\mathbf{X} = \mathbf{x}) = L_m - c_m \ell - \frac{d_m}{m} n.$$

274 Three things are equivalent: (i) the vector \mathbf{X} satisfies the identities (2.1); (ii) the pair
 275 (S_m, T_m) is equal to (ℓ, n) ; (iii) the partition $\lambda = (\lambda_1, \dots, \lambda_m)$ defined by $\lambda_j - \lambda_{j+1} =$
 276 X_j for $2 \leq j \leq m-1$, together with $\lambda_1 = \ell - X_0$ and $\lambda_m = X_m$, is a partition of n
 277 fitting inside a $m \times \ell$ rectangle. Thus,

$$278 \quad N_n(\ell, m) = \mathbb{P}_m[(S_m, T_m) = (\ell, n)] \exp \left(-L_m + c_m \ell + \frac{d_m}{m} n \right)$$

$$279 \quad (3.9) \quad = \mathbb{P}_m[(S_m, T_m) = (\ell, n)] \exp \left[m \left(-\frac{L_m}{m} + c_m A + d_m B \right) \right].$$

281 Comparing (3.7) to (3.9), the proof is completed by an application of the LCLT in
 282 Lemma 5. \square

283 Lemma 5 is stated for an arbitrary sequence of parameters p_0, \dots, p_m bounded
 284 away from 0 and 1, though we need it only for $p_j = 1 - e^{-c_m - d_m j/m}$. For a 2×2
 285 matrix M , denote by $M(s, t) := [s, t] M [s, t]^T$ the corresponding quadratic form.

286 LEMMA 5 (LCLT). Fix $0 < \delta < 1$ and let p_0, \dots, p_m be any real numbers in
 287 the interval $[\delta, 1 - \delta]$. Let $\{X_j\}$ be independent reduced geometrics with respective
 288 parameters $\{p_j\}$, $S_m := \sum_{j=0}^m X_j$, and $T_m := \sum_{j=0}^m jX_j$. Let M_m be the covariance
 289 matrix for (S_m, T_m) , written

$$290 \quad M_m = \begin{pmatrix} \alpha_m m & \beta_m m^2 \\ \beta_m m^2 & \gamma_m m^3 \end{pmatrix},$$

291 Q_m denote the inverse matrix to M_m , and $\Delta_m = m^{-4} \det M_m = \alpha_m \gamma_m - \beta_m^2$. Let μ_m
 292 and ν_m denote the respective means $\mathbb{E}S_m$ and $\mathbb{E}T_m$. Denote $p_m(a, b) := \mathbb{P}((S_m, T_m) =$
 293 $(a, b))$. Then

$$294 \quad (3.10) \quad \sup_{a, b \in \mathbb{Z}} m^2 \left| p_m(a, b) - \frac{1}{2\pi(\det M_m)^{1/2}} e^{-\frac{1}{2}Q_m(a-\mu_m, b-\nu_m)} \right| \rightarrow 0$$

295 as $m \rightarrow \infty$, uniformly in the parameters $\{p_j\}$ in the allowed range. In particular, if
 296 the sequence (a_m, b_m) satisfies $Q_m(a_m - \mu_m, b_m - \nu_m) \rightarrow 0$ then

$$297 \quad \mathbb{P}(S_m = a_m, T_m = b_m) = \frac{1}{2\pi\sqrt{\Delta_m}m^2} \left(1 + O\left(m^{-3/2}\right) \right).$$

298 The following consequence will be used to prove Theorem 2.

299 COROLLARY 6 (LCLT consecutive differences). Define the normal approxima-
 300 tion $\mathcal{N}_m(a, b) := \frac{1}{2\pi(\det M_m)^{1/2}} e^{-\frac{1}{2}Q_m(a-\mu_m, b-\nu_m)}$ as in Equation (3.10). Using the
 301 notation of Lemma 5,

$$302 \quad \sup_{a, b \in \mathbb{Z}} \left| p_m(a, b+1) - p_m(a, b) - (\mathcal{N}_m(a, b+1) - \mathcal{N}_m(a, b)) \right| = O(m^{-4}).$$

303 The technical but unsurprising proofs of Lemma 5 and Corollary 6 are given in
 304 the Appendix at the end of this article.

305 **4. Limit shape.** Suppose a Young diagram is chosen uniformly from among
 306 all partitions of n fitting in a $m \times \ell$ rectangle. To simplify calculations, we imagine
 307 this Young diagram outlining a compact set in the fourth quadrant of the plane
 308 and rotate 90° counterclockwise to obtain a shape in the first quadrant. Let $\Xi_{n, m, \ell}$
 309 denote the random set obtained in this manner after rescaling by a factor of $1/m$, so
 310 that the length in the positive x -direction is bounded by 1. Fix $A > 2B > 0$ and
 311 metrize compact sets of \mathbb{R}^2 by the Hausdorff metric. As $m \rightarrow \infty$ with $\ell/m \rightarrow A$ and
 312 $n/m^2 \rightarrow B$, the random set $\Xi_{n, m, \ell}$ converges in distribution to a deterministic set
 313 $\Xi^{A, B}$. See Figure 4 for some examples.

314 Our methods immediately recover the distributional convergence result $\Xi_{n, m, \ell} \rightarrow$
 315 $\Xi^{A, B}$. As previously mentioned, this limit shape was known to Petrov [17] and others.
 316 Petrov identifies it as a portion of the limit curve for unrestricted partitions, which
 317 itself was posed as a problem by Vershik and answered in [27, 28] (see also [33]).
 318 Because this result is already known, along with precise fluctuation information which
 319 we do not derive, we give only the short argument here for distributional convergence.
 320 We do not determine the best possible fluctuation results following from this method.

321 The shape $\Xi_{n,m,\ell}$ is determined by its boundary, a polygonal path obtained
 322 from a partition λ by filling in unit vertical connecting lines in the step function
 323 $x \mapsto m^{-1}\lambda_{\lfloor mx \rfloor}$. Recall that the probability measure \mathbb{P}_m restricted to the event
 324 $\{(S_m, T_m) = (\ell, n)\}$ gives all partitions counted by $N_n(m, \ell)$ equal probability and
 325 that \mathbb{P}_m gives the event $\{(S_m, T_m) = (\ell, n)\}$ probability $\Theta(m^{-2})$. Distributional con-
 326 vergence of $\Xi_{n,m,\ell}$ to $\Xi^{A,B}$ then follows from the following.

PROPOSITION 7. Fix $A > 2B > 0$. Define the maximum discrepancy by

$$\mathcal{M} := \max_{0 \leq j \leq m} \left| \sum_{i=0}^j \left(X_i - \frac{q_i}{p_i} \right) \right|.$$

Then for any $\varepsilon > 0$,

$$\mathbb{P}_m[\mathcal{M} \geq \varepsilon m] = o(m^{-2})$$

327 as $m \rightarrow \infty$ with $\ell/m \rightarrow A$ and $n/m \rightarrow B$.

Proof. This is a routine application of exponential moment bounds. By our definition of p_i , in this regime there exists $\delta > 0$ such that $p_i \in [\delta, 1 - \delta]$ for all i . Therefore, there are $\eta, K > 0$ such that for $\lambda < \eta$, the mean zero variables $X_i - q_i/p_i$ all satisfy $\mathbb{E} \exp(\lambda(X_i - q_i/p_i)) \leq \exp(K\lambda^2)$. Independence of the family $\{X_i\}$ then gives

$$\mathbb{E} \exp \left[\lambda \sum_{i=0}^j (X_i - p_i/q_i) \right] \leq e^{Km\lambda^2}$$

for all $j \leq m$. By Markov's inequality,

$$\mathbb{P}(|X_i - p_i/q_i| \geq \varepsilon m) \leq e^{Km\lambda^2 - \lambda \varepsilon m}.$$

328 Fixing $\lambda = 1/(2K)$ shows that this probability is bounded above by $\exp(-m/(4K))$.
 329 Hence, $\mathbb{P}(\mathcal{M} \geq \varepsilon m) \leq me^{-m/(4K)} = o(m^{-2})$ as desired. \square

330 To see that Proposition 7 implies the limit shape statement, let $\lambda_i := \ell - (X_0 +$
 331 $\dots + X_{i-1})$ so that

$$332 \quad y^{(m)}(i) := \mathbb{E}_m \lambda_i = \ell - \sum_{j=0}^{i-1} q_j/p_j.$$

333 Proposition 7 shows the boundary of Ξ_m to be within $o(m)$ of the step function $y^{(m)}(\cdot)$
 334 except with probability $o(m^{-2})$. Since \mathbb{P}_m restricted to the event $\{(S_m, T_m) = (\ell, n)\}$
 335 gives all partitions counted by $N_n(m, \ell)$ equal probability and \mathbb{P}_m gives the event
 336 $\{(S_m, T_m) = (\ell, n)\}$ probability $\Theta(m^{-2})$, the conditional law $(\mathbb{P}_m | (S_m, T_m) = (\ell, n))$
 337 gives the event $\{\mathcal{M} > \varepsilon m\}$ probability $o(1)$ as $m \rightarrow \infty$ with $\ell/m \rightarrow A$ and $n/m \rightarrow B$.
 338 Thus, the boundary of Ξ_m converges in distribution to the limit

$$339 \quad (4.1) \quad y(x) := \lim_{m \rightarrow \infty} m^{-1} y^{(m)}(\lfloor mx \rfloor).$$

340 Figure 4 shows examples of two families of the limit curve as well as a plot of the
 341 limit curve against uniformly generated restricted partitions for several values of m
 342 in the range $[120, 300]$.

Substituting the definition of $y^{(m)}(i)$ into (4.1) and evaluating the limit as an integral gives

$$y(x) = A + x - \int_0^x \frac{1}{1 - e^{-c-dt}} dt = A + x - \frac{1}{d} \ln \left(\frac{e^{xd+c} - 1}{e^c - 1} \right).$$

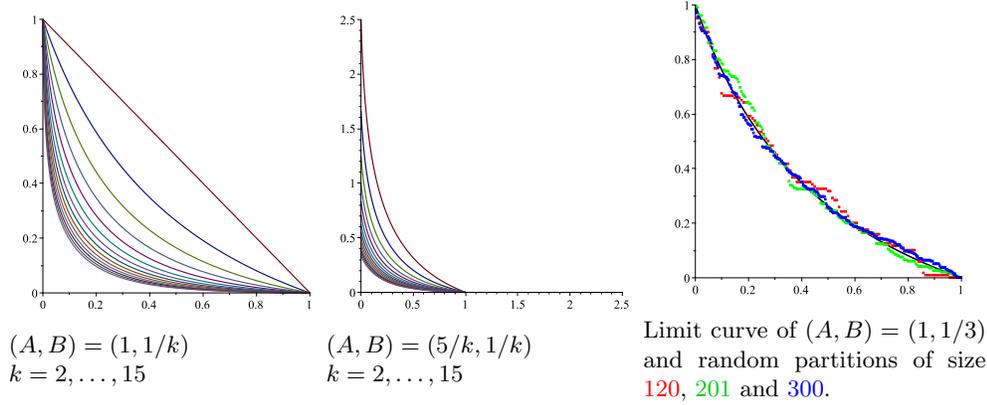


Fig. 4: Limit shapes of scaled partitions as $m \rightarrow \infty$.

After expressing c in terms of d , this may be written implicitly as

$$e^{(A+1)d} - 1 = (e^d - 1)e^{d(A-y)} + (e^{Ad} - 1)e^{d(1-x)}$$

343 which simplifies to

$$344 \quad (4.2) \quad (1 - e^{-c})e^{d(A-y)} + e^{-c}e^{-dx} = 1$$

345 as long as $A > 2B$; in the special case $A = 2B$ one obtains simply $y = A \cdot (1 - x)$.

346 It is worth comparing this result with the limit shape derived in [17]. There the
347 limit shape of the boxed partitions is identified as the portion of the curve $\{e^{-x} + e^{-y} =$
348 $1\}$, which is the limit shape of unrestricted partitions. The portion is determined
349 implicitly by the restriction that the endpoints of the curve are the opposite corners
350 of a $1 \times A$ -proportional rectangle and that the area under the curve has the desired
351 proportion, that is B/A of the total rectangular area. To see that this matches (4.2)
352 we can calculate the given portion explicitly.

353 Let $x = s_1, s_2$ be the starting and ending points of the bounding rectangle. The
354 side ratio and the area requirement are respectively equivalent to

$$355 \quad \frac{\log(1 - e^{s_1}) - \log(1 - e^{-s_2})}{s_2 - s_1} = A$$

356 and

$$357 \quad \int_{s_1}^{s_2} -\log(1 - e^{-t}) dt + (s_2 - s_1) \log(1 - e^{-s_2}) = B(s_2 - s_1)^2$$

358

359 which simplify to

$$360 \quad (4.3) \quad A = \frac{1}{s_2 - s_1} \log \left(\frac{e^{s_2} - 1}{e^{s_2} - e^{s_2 - s_1}} \right),$$

$$361 \quad (4.4) \quad B = \frac{-\operatorname{dilog}(1 - e^{-s_2}) + \operatorname{dilog}(1 - e^{-s_1}) + (s_2 - s_1) \log(1 - e^{-s_2})}{(s_2 - s_1)^2}.$$

Comparing these equations with equations (1.2) and (1.3) it is immediate that the solutions are given by $s_1 = c$ and $s_2 = c + d$. Finally, to match the curve in the

second line of equation (4.2) we need the coordinate transform from the curve γ in the segment $x = [c, c + d]$ given by

$$x \rightarrow x_1 = \frac{(x - c)}{d}, \quad y \rightarrow y_1 - A = \frac{y + \log(1 - e^{-c})}{d}$$

whence $x = dx_1 + c$ and $y = -d(A - y_1) - \log(1 - e^{-c})$ and the curves match.

5. Existence and Uniqueness of c, d . We now show that for any $A \geq 2B > 0$ there exists unique positive constants c and d satisfying Equations (1.2) and (1.3). If $A = B/2$ then $d = 0$ and c can be determined uniquely, so we may assume $A > 2B > 0$. The following lemma will be used to show uniqueness.

LEMMA 8. *Let ψ denote the map taking the pair (c, d) to (A, B) defined by the two integrals in Equations (1.2) and (1.3), and let K be a compact subset of $\{(x, y) : x > 2y > 0\}$. The Jacobian matrix $J := D[\psi]$ is negative definite for all $(c, d) \in (0, \infty)^2$, and all entries of ψ and J (respectively ψ^{-1} and J^{-1}) are Lipschitz continuous on $\psi^{-1}[K]$ (respectively K).*

Proof. Differentiating under the integral sign shows that the partial derivatives comprising the entries of $D[\psi]$ are given by

$$\begin{aligned} J_{A,c} &= \int_0^1 \frac{-e^{-(c+dt)}}{(1 - e^{-(c+dt)})^2} dt \\ J_{A,d} &= \int_0^1 \frac{-t e^{-(c+dt)}}{(1 - e^{-(c+dt)})^2} dt \\ J_{B,c} &= \int_0^1 \frac{-t e^{-(c+dt)}}{(1 - e^{-(c+dt)})^2} dt \\ J_{B,d} &= \int_0^1 \frac{-t^2 e^{-(c+dt)}}{(1 - e^{-(c+dt)})^2} dt; \end{aligned}$$

note that each term is negative. Let ρ denote the finite measure on $[0, 1]$ with density $e^{-(c+dt)}/(1 - e^{-(c+dt)})^2$ and let \mathbb{E}_ρ denote expectation with respect to ρ . Then

$$J_{A,c} = \mathbb{E}_\rho[-1], \quad J_{A,d} = J_{B,c} = \mathbb{E}_\rho[-t], \quad J_{B,d} = \mathbb{E}_\rho[-t^2],$$

and

$$\det J = \mathbb{E}_\rho[1] \cdot \mathbb{E}_\rho[t^2] - (\mathbb{E}_\rho[t])^2 = \mathbb{E}_\rho[1]^2 \cdot \text{Var}_\sigma[t],$$

where $\text{Var}_\sigma[t]$ denotes the variance of t with respect to the normalized measure $\sigma = \rho/\mathbb{E}_\rho[1]$. In particular, $\det J$ is positive, and bounded above and below when c and d are bounded away from 0, implying the stated results on Lipschitz continuity. As J is real and symmetric, it has real eigenvalues. Since the trace of J is negative while its determinant is positive, the eigenvalues of J have negative sum and positive product, meaning both are strictly negative and J is negative definite for any $c, d > 0$. \square

LEMMA 9. *For any $A > 0$ and $B \in (0, A/2)$ there exist unique $c, d > 0$ satisfying Equations (1.2) and (1.3). Moreover, for a fixed A , when B decreases from $A/2$ to 0 then d increases strictly from 0 to ∞ and c decreases strictly from $\log(\frac{A+1}{A})$ to 1. When $B > 0$ is fixed and A goes to ∞ then c goes to 0 and d goes to the root of*

$$d^2 = B(d \log(1 - e^{-d}) - \text{dilog}(1 - e^{-d})).$$

389 *Proof.* Solving Equation (1.2) for c (assuming $d \geq 0$) gives

$$390 \quad c = \log \left(\frac{e^{(A+1)d} - 1}{e^{(A+1)d} - e^d} \right).$$

391 Substituting this into Equation (1.3) gives an explicit expression for B in terms of A
 392 and d , and shows that for fixed $A > 0$ as d goes from 0 to infinity B goes from $A/2$ to
 393 0. By continuity, this implies the existence of the desired c and d . It also shows that,
 394 for a fixed A , c is a decreasing function of d with the given maximal and minimal
 395 values as d goes from 0 to ∞ .

396 To prove uniqueness, we note that for $\mathbf{x}, \mathbf{y} \in \mathbb{R}^2$ Stokes' theorem implies

$$397 \quad \psi(\mathbf{y}) - \psi(\mathbf{x}) = \int_0^1 D[\psi](t\mathbf{x} + (1-t)\mathbf{y}) \cdot (\mathbf{x} - \mathbf{y}) dt$$

398 so that

$$399 \quad (\mathbf{x} - \mathbf{y})^T \cdot (\psi(\mathbf{y}) - \psi(\mathbf{x})) = \int_0^1 [(\mathbf{x} - \mathbf{y})^T \cdot D[\psi](t\mathbf{x} + (1-t)\mathbf{y}) \cdot (\mathbf{x} - \mathbf{y})] dt.$$

400 When $\mathbf{x} \neq \mathbf{y}$, negative-definiteness of $D[\psi]$ implies that the last integrand is strictly
 401 negative on $[0, 1]$, and $\psi(\mathbf{y}) \neq \psi(\mathbf{x})$. Thus, distinct values of c and d give distinct
 402 values of A and B .

403 To see the monotonicity, let A be fixed and let $F_B(d) = B$ be the equation ob-
 404 tained after substituting $c = c(A, d)$ above in Equation (1.3), i.e. $F_B(d) = \psi_2(c(A, d), d)$.
 405 Then d is a decreasing function of B and vice versa since

$$406 \quad \frac{\partial F_B(d)}{\partial d} = \frac{J_{B,d} J_{A,c} - J_{A,d} J_{B,c}}{J_{A,c}} = \frac{\det D[\psi]}{J_{A,c}} < 0. \quad \square$$

407 For the last part, the explicit formula for c in terms of A and d shows that $c \rightarrow 0$.
 408 Substitution in Equation (1.3) gives the desired equation.

409 **6. Proof of Theorem 1 from the discretized result.** Here we show how c_m
 410 and d_m from the discretized result are related to c, d defined independently of m . The
 411 proof below also shows that c_m and d_m exist and are unique.

412 The Euler-MacLaurin summation formula [6, Section 3.6] gives an expansion

$$413 \quad \frac{L_m}{m} = \int_0^1 \log(1 - e^{-c_m - d_m t}) dt + \frac{\log(1 - e^{-c_m}) + \log(1 - e^{-c_m - d_m})}{2m} + O(m^{-2})$$

$$(6.1)$$

$$414 \quad = \frac{\operatorname{dilog}(1 - e^{-c_m - d_m}) - \operatorname{dilog}(1 - e^{-c_m})}{d_m} + \frac{\log(1 - e^{-c_m}) + \log(1 - e^{-c_m - d_m})}{2m} + O(m^{-2}) \quad \blacksquare$$

416 of the sum L_m in terms of c_m and d_m . Assume that there is an asymptotic expansion

$$417 \quad (6.2) \quad c_m = c + um^{-1} + O(m^{-2})$$

$$418 \quad (6.3) \quad d_m = d + vm^{-1} + O(m^{-2})$$

419 as $m \rightarrow \infty$, where u and v are constants depending only on A and B . Under such an
 420 assumption, substitution of Equations (6.2) and (6.3) into Equation (6.1) implies

$$421 \quad \frac{L_m}{m} = \frac{\operatorname{dilog}(1 - e^{-c-d}) - \operatorname{dilog}(1 - e^{-c})}{d} + \frac{uA + vB}{m} + O(m^{-2})$$

$$422 \quad (6.4) \quad = \log(1 - e^{-c-d}) - dB + \frac{uA + vB}{m} + O(m^{-2}).$$

423

424 Substituting Equations (6.2)–(6.4) into Equation (3.7) of Theorem 4 and taking the
425 limit as $m \rightarrow \infty$ then gives Theorem 1, as

$$426 \quad \Delta_m \rightarrow \left(\int_0^1 \frac{e^{-c-dt}}{(1-e^{-c-dt})^2} dt \right) \left(\int_0^1 \frac{t^2 e^{-c-dt}}{(1-e^{-c-dt})^2} dt \right) - \left(\int_0^1 \frac{te^{-c-dt}}{(1-e^{-c-dt})^2} dt \right)^2 = \Delta.$$

427

428 It remains to show the expansions in Equations (6.2) and (6.3). For $x, y > 0$,
429 define

$$430 \quad \bar{S}_m(x, y) := \frac{1}{m} \sum_{j=0}^m \frac{1}{1 - e^{-(x+yj/m)}} - 1,$$

$$431 \quad \bar{T}_m(x, y) := \frac{1}{m} \sum_{j=0}^m \frac{j/m}{1 - e^{-(x+yj/m)}} - \frac{1}{2}.$$

432

433 Another application of the Euler-MacLaurin summation formula implies

$$434 \quad (6.5) \quad \bar{S}_m(c, d) = A + A_1(c, d)m^{-1} + O(m^{-2}),$$

$$435 \quad (6.6) \quad \bar{T}_m(c, d) = B + B_1(c, d)m^{-1} + O(m^{-2}),$$

436 with

$$437 \quad A_1 = \frac{1}{2} \left(\frac{1}{1 - e^{-c}} + \frac{1}{1 - e^{-c-d}} \right) \quad \text{and} \quad B_1 = \frac{1}{2(1 - e^{-c-d})}.$$

438 Let \mathcal{J} denote the Jacobian $D[\psi]$ of the map ψ , introduced in Lemma 8, with respect
439 to c and d , and let

$$440 \quad (c'_m, d'_m) = (c, d) - m^{-1} \mathcal{J}^{-1} \cdot (A_1 - 1, B_1 - 1/2)^T.$$

441 A Taylor expansion around the point (c, d) gives

$$442 \quad \begin{aligned} 443 \quad \left(\begin{array}{c} \bar{S}_m(c'_m, d'_m) \\ \bar{T}_m(c'_m, d'_m) \end{array} \right) &= \left(\begin{array}{c} \bar{S}_m(c, d) \\ \bar{T}_m(c, d) \end{array} \right) - (\mathcal{J} + O(m^{-1})) \cdot \left(m^{-1} \mathcal{J}^{-1} \begin{pmatrix} A_1 \\ B_1 \end{pmatrix} \right) + O(m^{-2}) \\ 444 \quad &= \begin{pmatrix} A - 1/m \\ B - 1/2m \end{pmatrix} + O(m^{-2}) \\ 445 \quad &= \begin{pmatrix} \bar{S}_m(c_m, d_m) \\ \bar{T}_m(c_m, d_m) \end{pmatrix} + O(m^{-2}), \end{aligned}$$

446 where Equations (6.5) and (6.6) were used to approximate the Jacobian of ψ_m :
447 $(x, y) \mapsto (\bar{S}_m(x, y), \bar{T}_m(x, y))$ with respect to x and y .

448 The map ψ_m is Lipschitz for a similar reason as its continuous analogue. Namely,
449 consider the partial derivatives

$$450 \quad \begin{aligned} 451 \quad J_{S,x} &= \frac{1}{m} \sum_{j=0}^m -\frac{e^{-x-yj/m}}{(1 - e^{-x-yj/m})^2} \\ 452 \quad J_{S,y} &= \frac{1}{m^2} \sum_{j=0}^m -\frac{j e^{-x-yj/m}}{(1 - e^{-x-yj/m})^2} \\ 453 \quad J_{T,x} &= \frac{1}{m^2} \sum_{j=0}^m -\frac{j e^{-x-yj/m}}{(1 - e^{-x-yj/m})^2} \\ 454 \quad J_{T,y} &= \frac{1}{m^3} \sum_{j=0}^m -\frac{j^2 e^{-x-yj/m}}{(1 - e^{-x-yj/m})^2}. \end{aligned}$$

455 Let ρ_m be a discrete finite measure on $R_m := \{0, 1/m, 2/m, \dots, 1\}$ with density
 456 $e^{-x-yt}/(1 - e^{-x-yt})$ for $t \in R_m$ and 0 otherwise, and let \mathbb{E}_{ρ_m} be the expectation with
 457 respect to ρ_m . Then

$$458 \quad J_{S,x} = \mathbb{E}_{\rho_m}[-1], \quad J_{T,x} = J_{S,y} = \mathbb{E}_{\rho_m}[-t], \quad J_{T,y} = \mathbb{E}_{\rho_m}[-t^2]$$

459 and

$$460 \quad \det D[\psi_m] = \mathbb{E}_{\rho_m}[1]\mathbb{E}_{\rho_m}[t^2] - \mathbb{E}_{\rho_m}[t]^2 = \mathbb{E}_{\rho_m}[1]^2 \text{Var}_{\sigma_m}[t],$$

461 where σ_m is the probability function $\rho_m/\mathbb{E}_{\rho_m}[1]$. For any fixed m and (x, y) in a
 462 compact neighborhood of (A, B) , both the variance and the expectation are finite
 463 and bounded away from 0, as is the Jacobian determinant. Moreover, the trace
 464 $\text{Tr}D[\psi] = -\mathbb{E}_{\rho_m}[1 + t^2]$ is bounded away from 0 and infinity, so the Jacobian
 465 is negative definite with locally bounded eigenvalues, and hence ψ_m is locally Lipschitz.
 466 Since the norm of the Jacobian is bounded away from 0 and infinity, we have that the
 467 inverse map ψ_m^{-1} is also locally Lipschitz in a neighborhood of $\psi^{-1}(A, B)$. Moreover,
 468 similarly to proof of existence and uniqueness of c and d in Section 5, we have that
 469 there indeed are c_m and d_m as unique solutions of Equations (3.2) and (3.3) since the
 470 Jacobian is negative semi-definite.

471 The trapezoid formula implies $|J_{S,c} - J_{A,c}| = O(m^{-1})$, and similar bounds for the
 472 other differences of partial derivatives in the continuous and discrete settings. Hence,
 473 the bounds for the norms and eigenvalues of $D[\psi_m]$ are within $O(m^{-1})$ of the ones for
 474 $D[\psi]$, and ψ_m (and its inverse) is Lipschitz with a constant independent of m . Thus,

$$475 \quad O(m^{-2}) = \|\psi_m(c'_m, d'_m) - \psi_m(c_m, d_m)\| \geq C^{-1}\|(c'_m - c_m, d'_m - d_m)\|$$

476 for some constant C , so that the expansions (6.2) and (6.3) hold. \square

477 **7. Proof of Theorem 2.** We will prove Theorem 2 from Equation (3.9) and
 478 Corollary 6. Let $p_m(\ell, n) = \mathbb{P}_m[(S_m, T_m) = (\ell, n)]$ and let

$$479 \quad (7.1) \quad L_m(x, y) := \sum_{j=0}^m \log(1 - e^{-x-yj/m}),$$

$$480 \quad (7.2) \quad A_m(x, y) := \sum_{j=0}^m \frac{1}{1 - e^{-x-yj/m}} - (m+1),$$

$$481 \quad (7.3) \quad B_m(x, y) := \sum_{j=0}^m \frac{j/m}{1 - e^{-x-yj/m}} - \frac{m+1}{2}.$$

482
 483 Then c_m and d_m are the solutions to

$$484 \quad A_m(c_m, d_m) = \ell = Am, \quad B_m(c_m, d_m) = n/m = Bm.$$

485 Let c'_m, d'_m be the solutions to $A_m(c'_m, d'_m) = \ell$ and $B_m(c'_m, d'_m) = (n+1)/m$, and let
 486 $\Delta x = c'_m - c_m = O(m^{-2})$ and $\Delta y = d'_m - d_m = O(m^{-2})$ by the Lipschitz properties
 487 proven in Section 5. Observe that

$$488 \quad (7.4) \quad \frac{\partial L_m(x, y)}{\partial x} = A_m(x, y) \quad \text{and} \quad \frac{\partial L_m(x, y)}{\partial y} = B_m(x, y).$$

489 Using the Taylor expansion for $L_m(c'_m, d'_m)$ around (c_m, d_m) and the L_m partial de-
 490 rivatives from Equation (7.4),

$$491 \quad -L_m(c'_m, d'_m) = -L_m(c_m + \Delta x, d_m + \Delta y)$$

$$492 \quad = -L_m(c_m, d_m) - \Delta x A_m(c_m, d_m) - \Delta y B_m(c_m, d_m) + O(m^{-3}),$$

494 so that

$$496 \quad -L_m(c'_m, d'_m) + (c_m + \Delta x)\ell + (d_m + \Delta y)(n+1)m^{-1} = -L_m(c_m, d_m) + c_m\ell + d_m(n+1)m^{-1} + O(m^{-3}). \blacksquare$$

497 To lighten notation, we now write $L_m := L_m(c_m, d_m)$ and $L'_m := L_m(c'_m, d'_m)$. Then

$$498 \quad N_{n+1}(\ell, m) - N_n(\ell, m) = p_m(\ell, n+1) \exp \left[-L'_m + c'_m\ell + \frac{d'_m}{m}(n+1) \right] - p_m(\ell, n) \exp \left[-L_m + c_m\ell + \frac{d_m}{m}n \right]$$

(7.5)

$$499 \quad = p_m(\ell, n) \exp \left[-L_m + c_m\ell + \frac{d_m}{m}n \right] \left[e^{d_m/m} - 1 \right]$$

(7.6)

$$500 \quad + [p_m(\ell, n+1) - p_m(\ell, n)] \exp \left[-L_m + c_m\ell + \frac{d_m}{m}(n+1) \right]$$

(7.7)

$$501 \quad + p_m(\ell, n+1) \left(e^{-L'_m + c'_m\ell + d'_m(n+1)/m} - e^{-L_m + c_m\ell + d_m(n+1)/m} \right). \blacksquare$$

503 We now bound each of these summands.

- 504 • Since $d_m = d + O(m^{-1})$, Equation (3.9) implies that the quantity on line (7.5)
- 505 equals

$$506 \quad N_n(\ell, m) \left(\frac{d}{m} + O(m^{-2}) \right)$$

507 as long as $d \notin O(m^{-1})$. This holds when $|A - B/2| \notin O(m^{-1})$ as $d = 0$ when

508 $A = B/2$ and the map taking (A, B) to (c, d) is Lipschitz.

- 509 • By Corollary 6,

$$510 \quad [p_m(\ell, n+1) - p_m(\ell, n)] \leq |\mathcal{N}_m(\ell, n+1) - \mathcal{N}_m(\ell, n)| + O(m^{-4})$$

$$511 \quad = O \left(m^{-2} \cdot \left| 1 - e^{\frac{1}{2}Q_m(0,1)} \right| \right) + O(m^{-4})$$

$$512 \quad = O(m^{-4}),$$

514 where Q_m is the inverse of the covariance matrix of (S_m, T_m) . Thus, the

515 quantity on line (7.6) is $O(m^{-4} \cdot m^2 N_n(\ell, m)) = O(m^{-2} N_n(\ell, m))$.

- 516 • Let

$$517 \quad \psi_m := \exp \left[-L'_m + c'_m\ell + d'_m(n+1)m^{-1} - (-L_m + c_m\ell + d_m(n+1)m^{-1}) \right] - 1 = O(m^{-3}). \blacksquare$$

518 As $p_m(\ell, n+1) = p_m(\ell, n) + O(m^{-4})$, it follows that the quantity on line (7.7)

519 is

$$520 \quad p_m(\ell, n+1) e^{-L_m + c_m\ell + d_m(n+1)/m} \psi_m = N_n(\ell, m) \psi_m e^{d_m/m} + O(m^{-4} e^{d_m/m} e^{-L_m + c_m\ell + d_m n/m} \psi_m)$$

$$522 \quad = O(m^{-3} N_n(\ell, m)). \blacksquare$$

523 Putting everything together,

$$524 \quad N_{n+1}(\ell, m) - N_n(\ell, m) = N_n(\ell, m) \left(\frac{d}{m} + O(m^{-2}) \right),$$

525 as desired. \square

526 **Appendix: Proof of the Local Central Limit Theorem.** Throughout this
 527 section, $1/2 \geq \delta > 0$ is fixed and $\{p_j : 0 \leq j \leq m\}$ are arbitrary numbers in $[\delta, 1 - \delta]$.
 528 The variables $\{X_j\}$ and (S_m, T_m) are as in Lemma 5; we drop the index m on the
 529 remaining quantities $\alpha_m, \beta_m, \gamma_m, \Delta_m, \mu_m, \nu_m, p_m(a, b)$ and the matrices M_m and Q_m .
 530 Recall the quadratic form notation $M(s, t) := [s, t] M [s, t]^T$.

531 **LEMMA 10.** *The constants α, β, γ and Δ are bounded below and above by positive*
 532 *constants depending only on δ .*

533 **PROOF:** Upper and lower bounds on α, β and γ are elementary: $\alpha \in \left[\frac{\delta}{(1-\delta)^2}, \frac{(1-\delta)}{\delta^2} \right]$,

534 $\beta \in \left[\frac{\delta}{2(1-\delta)^2}, \frac{(1-\delta)}{2\delta^2} \right]$ and $\gamma \in \left[\frac{\delta}{3(1-\delta)^2}, \frac{(1-\delta)}{3\delta^2} \right]$. The upper bound on Δ follows
 535 from these.

536 For the lower bound on Δ , let $\tilde{M} = \begin{pmatrix} \alpha_n & \beta_n \\ \beta_n & \gamma_n \end{pmatrix}$ denote M without the factors of
 537 m . We show Δ is bounded from below by the positive constant $(4 - \sqrt{13})\delta/6$. A lower
 538 bound for the determinant Δ of \tilde{M} is $|\lambda|^2$ where λ is the least modulus eigenvalue of
 539 \tilde{M} ; note that $|\lambda|^2 = \inf_{\theta} \tilde{M}(\cos \theta, \sin \theta)$. We compute

$$\begin{aligned} 540 \quad \tilde{M}(\cos \theta, \sin \theta) &= m^{-1} \mathbb{E} (\cos \theta S + m^{-1} \sin \theta T)^2 \\ 541 \quad &\geq \delta m^{-1} \sum_{k=0}^m \left(\cos \theta + \frac{k}{m} \sin \theta \right)^2 \\ 542 \quad &> \delta \cdot \left(\cos^2 \theta + \cos \theta \sin \theta + \frac{1}{3} \sin^2 \theta \right). \\ 543 \end{aligned}$$

544 This is at least $\frac{4 - \sqrt{13}}{6} \delta$ for all θ , proving the lemma. \square

545 **LEMMA 11.** *Let X_p denote a reduced geometric with parameter p . For every $\delta \in$*
 546 *$(0, 1/2)$ there is a K such that simultaneously for all $p \in [\delta, 1 - \delta]$,*

$$547 \quad \left| \log \mathbb{E} \exp(i\lambda X_p) - \left(i \frac{q}{p} \lambda - \frac{q^2}{2p^2} \lambda^2 \right) \right| \leq K \lambda^3.$$

548 **PROOF:** For fixed p this is Taylor's remainder theorem together with the fact that
 549 the characteristic function $\phi_p(\lambda)$ of X_p is thrice differentiable. The constant $K(p)$ one
 550 obtains this way is continuous in p on the interval $(0, 1)$, therefore bounded on any
 551 compact sub-interval. \square

552 **PROOF OF THE LCLT:** The proof of Lemma 5 comes from expressing the probabil-
 553 ity as an integral of the characteristic function, via the inversion formula, and then
 554 estimating the integrand in various regions.

555 Let $\phi(s, t) := \mathbb{E} e^{i(sS + tT)}$ denote the characteristic function of (S, T) . Centering
 556 the variables at their means, denote $\hat{S} := S - \mu$, $\hat{T} := T - \nu$, and $\hat{\phi}(s, t) := \mathbb{E} e^{i(s\hat{S} + t\hat{T})}$
 557 so that $\phi(s, t) = \hat{\phi}(s, t) e^{is\mu + it\nu}$. Then

$$\begin{aligned} 558 \quad p(a, b) &= \frac{1}{(2\pi)^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} e^{-isa - itb} \phi(s, t) ds dt \\ 559 \quad (7.8) \quad &= \frac{1}{(2\pi)^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} e^{-is(a-\mu) - it(b-\nu)} \hat{\phi}(s, t) ds dt. \\ 560 \end{aligned}$$

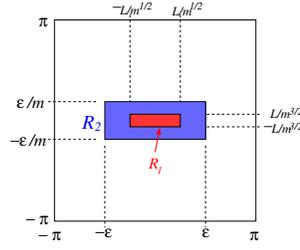


Fig. 5: The regions $\mathcal{R}_1 \subseteq \mathcal{R}_2 \subseteq \mathcal{R}$ in the proof of the LCLT.

561 Following the proof of the univariate LCLT for IID variables found in [7], we observe
 562 that
 (7.9)

$$563 \frac{1}{2\pi(\det M)^{1/2}} e^{-\frac{1}{2}Q(a-\mu, b-\nu)} = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-is(a-u)-it(b-v)} \exp\left(-\frac{1}{2}M(s,t)\right) ds dt.$$

564 Hence, comparing this to (7.8) and observing that $e^{-is(a-\mu)-it(b-\nu)}$ has unit modulus,
 565 the absolute difference between $p(a, b)$ and the left-hand side of (7.9) is bounded above
 566 by

$$567 (7.10) \quad \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left| \mathbf{1}_{(s,t) \in [-\pi, \pi]^2} \widehat{\phi}(s, t) - e^{-(1/2)M(s,t)} \right| ds dt.$$

568 Fix positive constants L and ε to be specified later and decompose the region
 569 $\mathcal{R} := [-\pi, \pi]^2$ as the disjoint union $\mathcal{R}_1 + \mathcal{R}_2 + \mathcal{R}_3$, where

$$570 \quad \mathcal{R}_1 = [-Lm^{-1/2}, Lm^{-1/2}] \times [-Lm^{-3/2}, Lm^{-3/2}]$$

$$571 \quad \mathcal{R}_2 = [-\varepsilon, \varepsilon] \times [-\varepsilon m^{-1}, \varepsilon m^{-1}] \setminus \mathcal{R}_1$$

$$572 \quad \mathcal{R}_3 = \mathcal{R} \setminus (\mathcal{R}_1 \cup \mathcal{R}_2);$$

574 see Figure 5 for details.

575 As $\int_{\mathcal{R}_2} e^{-(1/2)M(s,t)} ds dt$ decays exponentially with m , it suffices to obtain the
 576 following estimates

$$577 (7.11) \quad \int_{\mathcal{R}_1} \left| \widehat{\phi}(s, t) - e^{-(1/2)M(s,t)} \right| ds dt = O(m^{-5/2})$$

$$578 (7.12) \quad \int_{\mathcal{R}_2} \left| \widehat{\phi}(s, t) - e^{-(1/2)M(s,t)} \right| ds dt = O(m^{-5/2})$$

$$579 (7.13) \quad \int_{\mathcal{R}_3} \left| \widehat{\phi}(s, t) \right| ds dt = o(m^{-3}).$$

581 By independence of $\{X_j\}$,

$$582 \quad \log \widehat{\phi}(s, t) = \sum_{j=0}^m \log \mathbb{E} e^{i(s+jt)(X_j - \mu_j)}.$$

583 Using Lemma 11 with $p = p_j$ gives

$$584 \quad \left| \log \mathbb{E} e^{i(s+jt)(X_j - q_j/p_j)} + \frac{q_j}{2p_j^2} (s+jt)^2 \right| \leq K|s+jt|^3.$$

585 The sum of $(q_j/p_j^2)(s+jt)^2$ is $M(s,t)$, therefore summing the previous inequalities
586 over j gives

$$587 \quad (7.14) \quad \left| \log \widehat{\phi}(s,t) + \frac{1}{2}M(s,t) \right| \leq K \sum_{j=0}^m |s+jt|^3.$$

588 On \mathcal{R}_1 we have the upper bound $|s+jt| \leq |s|+m|t| \leq 2Lm^{-1/2}$. Thus,

$$589 \quad \sum_{j=0}^m |s+jt|^3 \leq (m+1)(8L^3)m^{-3/2} = O\left(m^{-1/2}\right).$$

590 Plugging this into (7.14) and exponentiating shows that the left hand side of (7.11)
591 is at most $|\mathcal{R}_1| \cdot O(m^{-1/2}) = O(m^{-5/2})$.

592 To bound the integral on \mathcal{R}_2 , we define the sub-regions

$$593 \quad S_k := \left\{ (x,y) : k \leq \max\left(m^{1/2}|x|, m^{3/2}|y|\right) \leq k+1 \right\}.$$

594 As the area of S_k is $(8k+4)m^{-2}$,

$$595 \quad \int_{\mathcal{R}_2} \left| \widehat{\phi}(s,t) - e^{-(1/2)M(s,t)} \right| ds dt \leq \sum_{k=L}^{\lceil \epsilon\sqrt{m} \rceil} \int_{S_k} \left| \widehat{\phi}(s,t) - e^{-M(s,t)/2} \right| ds dt$$

$$596 \quad (7.15) \quad \leq m^{-2} \sum_{k=L}^{\lceil \epsilon\sqrt{m} \rceil} (8k+4) \max_{(s,t) \in S_k} \left| \widehat{\phi}(s,t) - e^{-M(s,t)/2} \right|.$$

597

598 We break this last sum into two parts, and bound each part. For $(s,t) \in \mathcal{R}_2$, we have
599 $|s+jt| \leq |s|+m|t| \leq 2\epsilon$ so that

$$600 \quad \sum_{j=0}^m |s+jt|^3 \leq 2\epsilon \sum_{j=0}^m (|s|+j|t|)^2 \leq (2\epsilon\Delta^{-1})M(|s|,|t|).$$

601 Comparing this to (7.14) shows we may choose ϵ small enough to guarantee that

$$602 \quad \left| \log \widehat{\phi}(s,t) + \frac{1}{2}M(s,t) \right| \leq \frac{1}{4}M(|s|,|t|),$$

603 so $|\widehat{\phi}(s,t)| \leq e^{-(1/4)M(s,t)}$. Lemma 10 shows there is a positive constant c such that
604 the minimum value of $M(s,t)$ on S_k is at least ck^2 . Thus, for $(s,t) \in S_k$,

$$605 \quad \left| \widehat{\phi}(s,t) - e^{-M(s,t)/2} \right| \leq \left| e^{-M(s,t)/4} \right| + \left| e^{-M(s,t)/2} \right| \leq 2e^{-ck^2}.$$

606 If $r_m := \lceil \sqrt{(\log m)/c} \rceil$ then

$$607 \quad \sum_{k=r_m}^{\infty} (8k+4)(k+1) \max_{(s,t) \in S_k} \left| \widehat{\phi}(s,t) - e^{-M(s,t)/2} \right| \leq 2 \sum_{k=r_m}^{\infty} (8k+4)(k+1)e^{-ck^2}$$

$$608 \quad = O(m^{-1} \text{polylog}(m))$$

$$609 \quad (7.16) \quad = O(m^{-1/2}),$$

610

611 where $\text{polylog}(m)$ denotes a quantity growing as an integer power of $\log m$. Further-
 612 more, for $(s, t) \in S_k$ there exist constants C and C' such that

$$613 \quad \left| \log \widehat{\phi}(s, t) + M(s, t)/2 \right| \leq C \sum_{j=0}^m |s + jt|^3 \leq C \left(2(k+1)m^{-1/2} \right)^3 (m+1) = C' k^3 m^{-1/2}.$$

614 This implies the existence of a constant $K > 0$ such that for $0 \leq k \leq r_m$ and
 615 $(s, t) \in S_k$,

$$616 \quad \left| \widehat{\phi}(s, t) - e^{-M(s, t)/2} \right| = \left| e^{-M(s, t)/2} \left| 1 - e^{\log \widehat{\phi}(s, t) + M(s, t)/2} \right| \right| \\ 617 \quad \leq K e^{-ck^2} k^3 m^{-1/2}.$$

619 Thus,

$$620 \quad \sum_{k=L}^{r_m} (8k+4)(k+1) \max_{(s, t) \in S_k} \left| \widehat{\phi}(s, t) - e^{-M(s, t)/2} \right| \leq K m^{-1/2} \sum_{k=L}^{r_m} (8k+4)(k+1) k^3 e^{-ck^2} \\ 622 \quad (7.17) \quad = O(m^{-1/2}).$$

623 Combining (7.15)–(7.17) gives (7.12).

624 Finally, for (7.13), we claim there is a positive constant c for which $|\widehat{\phi}(s, t)| \leq e^{-cm}$
 625 on \mathcal{R}_3 . To see this, observe (see [7, p. 144]) that for each p there is an $\eta > 0$ such
 626 that $|\phi_p(\lambda)| < 1 - \eta$ on $[-\pi, \pi] \setminus [-\varepsilon/2, \varepsilon/2]$. Again, by continuity, we may choose one
 627 such η valid for all $p \in [\delta, 1 - \delta]$. It suffices to show that when either $|s|$ or $m|t|$ is at
 628 least ε , then at least $m/3$ of the summands $\log \mathbb{E} e^{i(s+jt)(X_j - \mu_j)}$ have real part at most
 629 $-\eta$. Suppose $s \geq \varepsilon$ (the argument is the same for $s \leq -\varepsilon$). Interpreting $s + jt$ modulo
 630 2π always to lie in $[-\pi, \pi]$, the number of $j \in [0, m]$ for which $s + jt \in [-\varepsilon/2, \varepsilon/2]$ is
 631 at most twice the number for which $s + jt \in [\varepsilon/2, \varepsilon]$, hence at most twice the number
 632 for which $s + jt \notin [-\varepsilon/2, \varepsilon/2]$; thus at least $m/3$ of the $m + 1$ values of $s + jt$ lie
 633 outside $[-\varepsilon/2, \varepsilon/2]$ and these have real part of $\log \mathbb{E} e^{i(s+jt)(X_j - \mu_j)} \leq -\eta$ by choice of
 634 η . Lastly, if instead one assumes $\pi \geq t \geq \varepsilon/m$, then at most half of the values of
 635 $s + jt$ modulo 2π can fall inside any interval of length $\varepsilon/2$. Choosing η such that the
 636 real part of $\log \mathbb{E} e^{i(s+jt)(X_j - \mu_j)}$ is at most $-\eta$ outside of $[-\varepsilon/4, \varepsilon/4]$ finishes the proof
 637 of (7.13) and the LCLT. \square

638 **PROOF OF COROLLARY 6.** In order to estimate the error terms in the approxima-
 639 tion of $p(a, b)$ we will consider the partial differences and repeat the approximation
 640 arguments above. Changing b to $b + 1$ in Equations (7.8) and (7.9) implies
 (7.18)

$$641 \quad \left| p(a, b+1) - p(a, b) - (\mathcal{N}(a, b+1) - \mathcal{N}(a, b)) \right| = \int_{[-\pi, \pi]^2} \left| 1 - e^{-it} \right| \left| \widehat{\phi}(s, t) - e^{-1/2M(s, t)} \right| ds dt.$$

642 For $(s, t) \in \mathcal{R}_3$, the proof of the LCLT shows that the integral in Equation (7.18)
 643 decays exponentially with m . As $|1 - e^{-it}| = \sqrt{2 - 2\cos(t)} \leq |t| = O(m^{-3/2})$ for
 644 $(s, t) \in \mathcal{R}_1$, the proof of the LCLT shows that the integral in Equation (7.18) grows as
 645 $O(m^{-3/2} \cdot m^{-5/2}) = O(m^{-4})$. Finally, since $|1 - e^{-it}| \leq |t| \leq (k+1)m^{-3/2}$ for $(s, t) \in$
 646 S_k following the proof of the LCLT shows $\int_{\mathcal{R}_2} |1 - e^{-it}| \left| \widehat{\phi}(s, t) - e^{-1/2M(s, t)} \right| ds dt$ is
 647 at most

$$649 \quad m^{-7/2} \sum_{k=L}^{\lceil \varepsilon\sqrt{m} \rceil} (8k+4)(k+1) \max_{(s, t) \in S_k} \left| \widehat{\phi}(s, t) - e^{-M(s, t)/2} \right| = O(m^{-4}).$$

651 \square

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 654 statistical mechanical ideas drawn on in this paper are already present.

655

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