

# Quantum Graphs and Quantum Cuntz-Krieger Algebras

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# Today's Goals

- ▶ To describe a generalization of **graph theory** within the framework of operator algebras/non-commutative geometry: [Quantum Graphs](#).
- ▶ To explain (briefly) why one might want to study quantum graphs.
- ▶ To introduce a new class of operator algebras that encode the “**symbolic dynamics**” associated to quantum graphs: [Quantum Cuntz-Krieger algebras](#).

## Joint work with:

- ▶ Kari Eifler (TAMU)
- ▶ Christian Voigt (Glasgow)
- ▶ Moritz Weber (Saarbrücken)

# Graphs

- A (finite, directed) **graph** is a tuple  $\mathcal{G} = (V, E, s, t)$  where
- ▶  $V, E$  are finite sets (**vertices** and **edges**, respectively)
  - ▶  $s, t : E \rightarrow V$  are functions (**source** and **target** maps).

# Graphs

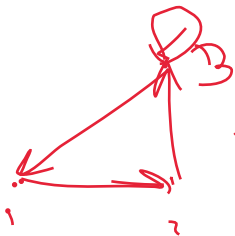
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  - ▶  $s, t : E \rightarrow V$  are functions (**source** and **target** maps).
  - ▶ Thus each edge  $e \in E$  is thus associated to an ordered pair  $(s(e), t(e)) \in V \times V$ .  $s(e)$  is the **source of  $e$**  and  $t(e)$  is the **target of  $e$** .
  - ▶ Graphs  $\mathcal{G}$  can **multiple edges**: The map  $E \rightarrow V \times V; e \mapsto (s(e), t(e))$  need not be injective.

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  - ▶ Graphs  $\mathcal{G}$  can **multiple edges**: The map  $E \rightarrow V \times V; e \mapsto (s(e), t(e))$  need not be injective.
  - ▶ For our purposes, we will **ONLY** consider graphs  $\mathcal{G}$  with **no multiple edges**.
  - ▶ In this case, we simply identify  $E \subseteq V \times V$ . In this case we can consider the **adjacency matrix**  $A_G \in M_{V \times V}(\{0, 1\})$  of  $\mathcal{G}$ :

$$A_G(v, w) = 1 \iff (v, w) \in E.$$

Then  $\mathcal{G}$  is encoded by the pair  $(V, A_G)$ .



$$A_G = \begin{matrix} & \begin{matrix} 1 & 2 & 3 \end{matrix} \\ \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} & \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 1 \end{pmatrix} \end{matrix}$$

# How to “Quantize” a Graph: The NCG Perspective

Non-commutative Geometry (NCG) is all about finding **quantum** (i.e., non-commutative) generalizations of familiar topological/algebraic/geometric structures by generalizing **Gelfand duality**:

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**Examples:**

- ▶ **LC Hausdorff spaces**  $X \longleftrightarrow$  abelian  $C^*$ -algebras  
 $A = C_0(X)$ .
- ▶ **Quantum Topological Spaces**  $\longleftrightarrow$  arbitrary  $C^*$ -algebras  
 $A \subset B(H)$ .
- ▶ **Measure spaces**  $(X, \mu) \longleftrightarrow$  abelian von Neumann algebras  
 $M = L^\infty(X, \mu)$ .
- ▶ **Quantum Measure spaces**  $\longleftrightarrow$  arbitrary von Neumann algebras  $M = M'' \subset B(H)$ .
- ▶ **Quantum groups**  $\longleftrightarrow$  nice Hopf algebras/tensor categories.
- ▶ **Quantum Riemannian manifolds**  $\longleftrightarrow$  Connes' spectral triples.



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- ▶ **Quantum groups**  $\longleftrightarrow$  nice Hopf algebras/tensor categories.
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**Our Problem:** **What is the quantum analogue of a graph?**

## Quantizing Graphs

**Input:** A finite directed graph  $\mathcal{G} = (V, A_{\mathcal{G}})$  with  $|V| = n$  vertices and no multiple edges.

**Output:** A  $C^*$ -algebraic description of  $\mathcal{G}$  that can be made non-commutative!

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- ▶ Let  $B = C(V) = C^*((p_v)_{v \in V} \mid p_v^* = p_v^2 = p_v, \sum_v p_v = 1)$ , the  $C^*$ -algebra of functions on  $V$ .
- ▶ Equip  $B$  with the canonical trace functional  $\psi : B \rightarrow \mathbb{C}$  given by counting measure on  $V$ .  $\psi$  is “**canonical**” because

$$\psi = \text{Tr}_{\text{End}(B)} \Big|_B \quad \text{where } B \hookrightarrow \text{End}(B) = \text{left-regular representation.}$$

- ▶ Use (GNS) inner product  $\langle f | g \rangle = \psi(f^* g)$  to turn  $B$  into a Hilbert space  $B = L^2(B) = \ell^2(V)$ .

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- ▶ Use (GNS) inner product  $\langle f | g \rangle = \psi(f^* g)$  to turn  $B$  into a Hilbert space  $B = L^2(B) = \ell^2(V)$ .
- ▶ View the adjacency matrix  $A_{\mathcal{G}}$  as a **linear map**  $A_{\mathcal{G}} \in \text{End}(B)$  in the obvious way:

$$A_{\mathcal{G}} p_w = \sum_{v \in V} A_{\mathcal{G}}(v, w) p_v.$$

## Quantizing Adjacency Matrices

So far:

$$\mathcal{G} \rightsquigarrow \left( B = C(V), \psi = \text{Tr}_{\text{End}(B)} \Big|_B, A_{\mathcal{G}} \in \text{End}(B) \right).$$

**Problem:** How to intrinsically capture the fact that  $A_{\mathcal{G}}$  is an **adjacency matrix** at the level of the function algebra  $B$ ?

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**Problem:** How to intrinsically capture the fact that  $A_{\mathcal{G}}$  is an **adjacency matrix** at the level of the function algebra  $B$ ?

**Answer:**  $A_{\mathcal{G}} \in M_n(\{0, 1\}) \iff A_{\mathcal{G}}$  is **idempotent** with respect to **Schur multiplication** in  $\text{End}(B)$ !

- ▶ Let  $m : B \otimes B \rightarrow B$  be the algebra multiplication,  $m^* : B \rightarrow B \otimes B$  its Hilbertian adjoint.

$$m(p_v \otimes p_w) = \delta_{v,w} p_v, \quad m^*(p_v) = (p_v \otimes p_v).$$

- ▶ Simple calculation: Given  $X, Y \in \text{End}(B)$ ,

$$\text{Schur Multiplication} : X \star Y = m(X \otimes Y)m^*$$

$$\text{Conclusion: } A_{\mathcal{G}} \in M_n(\{0, 1\}) \iff m(A_{\mathcal{G}} \otimes A_{\mathcal{G}})m^* = A_{\mathcal{G}}.$$

# Quantum Graphs

**We now go non-commutative!** Let  $B$  be a finite-dimensional  $C^*$ -algebra equipped with its **canonical trace functional**

$$\psi : B \rightarrow \mathbb{C}; \quad \psi = \text{Tr}_{\text{End}(B)} \Big|_B$$

(If  $B = \bigoplus_k M_{n(k)}(\mathbb{C})$ ),  $\psi = \sum_k n(k) \text{Tr}_{n(k)}(\cdot) =$  **Plancharel trace**).

## Definition

A linear map  $A \in \text{End}(B)$  is called a **quantum adjacency matrix** if

$$m(A \otimes A)m^* = A \quad (A \text{ is a “} \mathbf{\text{quantum Schur idempotent}} \text{”})$$

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We call the triple  $\mathcal{G} = (B, \psi, A)$  a **quantum graph**.

- ▶ Classical graphs  $\iff$  quantum graphs with **abelian**  $B$ .
- ▶ For certain applications, we can consider quantum graphs  $\mathcal{G} = (B, \psi, A)$  with **non-tracial states**  $\psi$ ...but not today.



## Some Basic Examples

Fix any pair  $(B, \psi)$ .

1. The **complete quantum graph** over  $B$  is  $\mathcal{K}(B) = (B, \psi, A)$ ,

$$A = \psi(\cdot)1_B \quad \text{Classical case: } A = \begin{pmatrix} 1 & \cdots & 1 \\ \vdots & 1 & \vdots \\ 1 & \cdots & 1 \end{pmatrix}$$

2. The **trivial quantum graph** over  $B$  is  $\mathcal{T}(B) = (B, \psi, A)$  where

$$A = \text{id}_{B \rightarrow B} \quad \text{Classical case: } A = \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & 1 & \vdots \\ 0 & \cdots & 1 \end{pmatrix}$$

3. If  $B = \bigoplus_k M_{n(k)}(\mathbb{C})$  and  $d = (d^{(k)})_k \in \bigoplus_k D_{n(k)} \subset B$  is **diagonal** with  $\text{Tr}_{n(k)}(d^{(k)}) = 1$  for all  $k$ , then

$$A_d(x) = dx \quad \text{is a quantum adjacency matrix.}$$

We call  $\mathcal{D}(B) = (B, \psi, A_d)$  a **diagonal quantum graph** over  $B$ .

## More Examples: Quantum Edge Space Picture

Writing down interesting examples of quantum adjacency matrices might at first seem not so easy.

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**Alternate Approach [Nik Weaver]:** Recall that a graph  $\mathcal{G}$  is completely described by its **edge relation**  $E \subset V \times V$ .

- ▶ Consider the linear subspace

$$S_{\mathcal{G}} = \text{span}\{e_{x,y} : (x,y) \in E\} \subset B(\ell^2(V)).$$

- ▶ Then  $S_{\mathcal{G}}$  is a **bimodule** over  $C(V) = C(V)' \subset B(\ell^2(V))$ .
- ▶ (Weaver) Every  $C(V)' - C(V)'$ -bimodule  $S \subset B(\ell^2(V))$  is of the form  $S_{\mathcal{G}}$  for some unique  $\mathcal{G} = (V, E)$ .

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### Definition (Weaver 2010)

Let  $B \subseteq B(H)$  be a finite dimensional  $C^*$ -algebra. A **quantum graph** on  $B$  is  $B'-B'$ -bimodule

$${}_{B'}S_{B'} \subseteq B(H).$$

Think of  ${}_{B'}S_{B'}$  as the collection of all “edge operators” that “connect quantum vertices” in the graph.

# Quantum Edge Space vs Quantum Adjacency Matrices

Both approaches to quantum graphs are essentially the same:

- ▶ **Fact 1:** Given  $B \subseteq B(H)$ , a bimodule  ${}_B S_{B'} \subseteq B(H)$  is uniquely determined by a self-adjoint projection

$$p = \sum_i a_i \otimes b_i^{op} \in B \otimes B^{op}; S = p \cdot B(H) = \left\{ \sum_i a_i X b_i : X \in B(H) \right\}$$

- ▶ **Fact 2:** Any such  $p = p^* = p^2 \in B \otimes B^{op}$  arises as the “**Choi-Jamiołkowski**” matrix of a **completely positive** quantum adjacency matrix  $A \in \text{End}(B)$  given by

$$p = (A \otimes 1)m^*(1_B).$$

**Conclusion:** Quantum graphs  $\mathcal{G} = (B, \psi, A)$  (with completely positive  $A$ )  $\longleftrightarrow$  operator spaces  $S \subseteq B(H)$  that are  $B'$ - $B'$ -bimodules.

## Quantum Graphs in “Nature”

Quantum graphs arise naturally in problems in quantum information theory and representation theory.

- ▶ **Zero-error channel capacity:** Given finite sets  $V, W$  and a **noisy channel** with transition probabilities  $(p(w|v))_{v \in V, w \in W}$ , can form the **confusability graph**  $\mathcal{G}_\Phi = (V, E)$  where

$$(v_1, v_2) \in E \iff \exists w \in W \text{ s.t. } p(w|v_1)p(w|v_2) \neq 0.$$

The **zero-error capacity of  $\Phi$** : Is the Shannon capacity of  $\mathbb{G}_\Phi$ .

- ▶ In **Quantum Information Theory**, one considers **quantum channels**

$$\Phi : L^1(B, \psi) \rightarrow M_k(\mathbb{C}).$$

The **Zero-error capacity of  $\Phi$** :  $\rightsquigarrow$  study Shannon capacity of the **non-commutative confusability graph  $S_\Phi$**  of  $\Phi$ :

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- ▶ Qgraphs  $\mathcal{G} = (B, \psi, A)$  appear in the theory of non-local games (graph isomorphism games), quantum teleportation and superdense coding schemes in QIT, representation theory of quantum symmetry groups of graphs, ...

## Cuntz-Krieger Algebras of graphs

Let  $\mathcal{G} = (V, A)$  be a finite graph without multiple edges.

The Cuntz-Krieger algebra  $\mathcal{O}(\mathcal{G}) = \mathcal{O}_A$  is the **universal  $C^*$ -algebra** generated by partial isometries  $(s_v)_{v \in V}$  satisfying the relations

$$(CK1) \quad \sum_{v \in V} s_v s_v^* = 1.$$

$$(CK2) \quad s_v^* s_v = \sum_{w \in V} A(v, w) s_w s_w^*.$$



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## Why study these operator algebras?

- ▶ These algebras turn out to encode the **symbolic dynamics** associated to a graph  $\mathcal{G}$ : They provide new conjugacy invariants for the **topological Markov chain**  $(X_A, \sigma_A)$  associated to  $\mathcal{G}$ .

$$X_A = \{(v_i)_{i \in \mathbb{N}} : A(v_i, v_{i+1}) = 1\} \text{ space of } \mathbf{infinite paths},$$
$$\sigma_A : X_A \rightarrow X_A; \quad \sigma_A(v_1, v_2, \dots) = (v_2, v_3, \dots).$$

- ▶ They provide a rich class of  $C^*$ -algebras that are amenable to classification.

## Examples

1.  $\mathcal{G} = K_n$  the complete graph on  $n$  vertices. Then  $\mathcal{O}(K_n)$  has generators  $(s_i)_{i=1}^n$  satisfying

$$(CK1) \sum_i s_i s_i^* = 1 \quad \& \quad (CK2) s_i^* s_i \left( = \sum_j s_j s_j^* \right) = 1.$$

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2.  $\mathcal{G} = T_n$  (trivial graph on  $n$  vertices, with self-loops). Then  $\mathcal{O}(T_n)$  has generators  $(s_i)_{i=1}^n$  satisfying

$$\sum_i s_i s_i^* = 1, \quad s_i^* s_i = s_i s_i^* \implies \mathcal{O}(T_n) = C(\mathbb{T})^{\oplus n}.$$

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$$\sum_i s_i s_i^* = 1, \quad s_i^* s_i = s_i s_i^* \implies \mathcal{O}(T_n) = C(\mathbb{T})^{\oplus n}.$$

3. Let  $\mathcal{G} \leftrightarrow A = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}$ . Then

$$\mathcal{O}(\mathcal{G}) = C^* \left( s_1, s_2 \mid \begin{array}{l} s_1 s_1^* + s_2 s_2^* = 1 \\ s_1^* s_1 = 1, s_2^* s_2 = s_1 s_1^* \end{array} \right).$$

## General properties of Cuntz-Krieger algebras

- ▶  $\mathcal{O}(\mathcal{G}) \neq 0$  as long as  $A_{\mathcal{G}} \neq 0$ .
- ▶  $\mathcal{O}(\mathcal{G})$  is always nuclear.
- ▶  $\mathcal{O}(\mathcal{G})$  is always unital.
- ▶  $\mathcal{O}(\mathcal{G})$  are classified, up to isomorphism.
- ▶ Include many naturally occurring examples of  $C^*$ -algebras.
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**Goal:** Generalize Cuntz-Krieger Algebras to Quantum Graphs!

## Quantum Cuntz-Krieger Algebras

Now let  $\mathcal{G} = (B, \psi, A)$  be a quantum graph. We want to build an operator algebra  $\mathbb{F}\mathcal{O}(\mathcal{G})$  out of  $\mathcal{G}$  as above. Here's our attempt.

## Quantum Cuntz-Krieger Algebras

Now let  $\mathcal{G} = (B, \psi, A)$  be a quantum graph. We want to build an operator algebra  $\mathbb{FO}(\mathcal{G})$  out of  $\mathcal{G}$  as above. Here's our attempt.

### Definition (B-Eifler-Voigt-Weber)

The **quantum Cuntz-Krieger Algebra** associated to  $\mathcal{G}$  is the universal  $C^*$ -algebra  $\mathbb{FO}(\mathcal{G})$  generated by the range of a linear map  $S : B \rightarrow \mathbb{FO}(\mathcal{G})$  satisfying the relations

1.  $\mu(1 \otimes \mu)(S \otimes S^* \otimes S)(1 \otimes m^*)m^* = S$
2.  $\mu(S^* \otimes S)m^* = \mu(S \otimes S^*)m^*A$

where  $\mu : \mathbb{FO}(\mathcal{G}) \otimes \mathbb{FO}(\mathcal{G}) \rightarrow \mathbb{FO}(\mathcal{G})$  is the multiplication map and  $S^*(b) = S(b^*)^*$ .



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## Why is this a reasonable definition?

### Theorem

*If  $\mathcal{G}$  is classical (i.e.,  $B = C(V)$  is abelian), then  $\mathbb{FO}(\mathcal{G})$  is the usual Cuntz-Krieger algebra  $\mathcal{O}(\mathcal{G})$  associated to  $\mathcal{G}$  (up to a  $KK$ -equivalence).*

## Quantum CK-Algebras: Unpacking the Definition

$$\mathcal{G} = (B, \psi, A) \rightsquigarrow \mathbb{FO}(\mathcal{G}) = C^* \left( (S(b))_{b \in B} \mid \begin{array}{l} S: B \rightarrow \mathcal{O}_{\mathcal{G}} \text{ linear} \\ \mu(1 \otimes \mu)(S \otimes S^* \otimes S)(1 \otimes m^*)m^* = S \\ \mu(S^* \otimes S)m^* = \mu(S \otimes S^*)m^* A \end{array} \right)$$

Write  $B = \bigoplus_{k=1}^m M_{n(k)}(\mathbb{C})$  with standard matrix unit basis  $\{e_{ij}^{(k)}\}$ .

For  $1 \leq k \leq m$ , put  $f_{ij}^{(k)} = n(k)^{-1} e_{ij}^{(k)}$  and

$$S^{(k)} = [S(f_{ij}^{(k)})] \in M_{n(k)}(\mathbb{FO}(\mathcal{G})).$$

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$$\mathcal{G} = (B, \psi, A) \rightsquigarrow \mathbb{FO}(\mathcal{G}) = C^* \left( (S(b))_{b \in B} \mid \begin{array}{l} S: B \rightarrow \mathcal{O}_{\mathcal{G}} \text{ linear} \\ \mu(1 \otimes \mu)(S \otimes S^* \otimes S)(1 \otimes m^*)m^* = S \\ \mu(S^* \otimes S)m^* = \mu(S \otimes S^*)m^* A \end{array} \right)$$

Write  $B = \bigoplus_{k=1}^m M_{n(k)}(\mathbb{C})$  with standard matrix unit basis  $\{e_{ij}^{(k)}\}$ .

For  $1 \leq k \leq m$ , put  $f_{ij}^{(k)} = n(k)^{-1}e_{ij}^{(k)}$  and

$$S^{(k)} = [S(f_{ij}^{(k)})] \in M_{n(k)}(\mathbb{FO}(\mathcal{G})).$$

Then  $\mu(1 \otimes \mu)(S \otimes S^* \otimes S)(1 \otimes m^*)m^* = S$

$$\iff S^{(k)} S^{(k)*} S^{(k)} = S^{(k)}, \quad k = 1, \dots, m$$

$\iff$  each  $S^{(k)}$  is a **partial isometry** in  $M_{n(k)}(\mathbb{FO}(\mathcal{G}))$ .

$$\mu(S^* \otimes S)m^* = \mu(S \otimes S^*)m^* A$$

$$\iff S^{(k)*} S^{(k)} = (S^{(1)} S^{(1)*}, \dots, S^{(m)} S^{(m)*}) \tilde{A}$$

where  $\tilde{A}_{ijk}^{xyz} = \frac{n(z)}{n(k)} A_{ijk}^{xyz}$  (reweighted adjacency matrix).

# Quantum CK-Algebras: Unpacking the Definition

**Summary:** Given  $\mathcal{G} = (B, \psi, A)$ ,  $B = \bigoplus_{k=1}^m M_{n(k)}$ , the quantum Cuntz-Krieger Algebra  $\mathcal{O}_{\mathcal{G}}$  is:

1. Generated by the coefficients of  $m$  **matrix partial isometries**  $S^{(k)} \in M_{n(k)}(\mathbb{F}\mathcal{O}(\mathcal{G}))$ ,
2. Subject to the **quantum Cuntz-Krieger relations**

$$S^{(k)*} S^{(k)} = (S^{(1)} S^{(1)*}, \dots, S^{(m)} S^{(m)*}) \tilde{A}.$$

3. **BUT:** There seems to be no natural analogue of the relation

$$\text{“} \sum_k S^{(k)} (S^{(k)})^* = 1 \text{”} \quad (\text{incompatible matrix sizes})$$

$\implies \mathcal{O}(\mathcal{G})$  might not be unital in general.

## Basic Examples, Basic Properties

**Example 1:** Trivial Quantum Graph  $\mathcal{G} = \mathcal{T}(M_n)$ .

$$\mathbb{FO}(\mathcal{T}(M_n)) = C^* \left( S_{ij}, 1 \leq i, j \leq n \mid \begin{array}{l} S = [S_{ij}] \text{ is a partial isometry} \\ \sum_l S_{li}^* S_{lj} = \sum_l S_{il} S_{jl}^* \iff S^* S = S S^* \end{array} \right)$$

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- ▶  $\mathbb{FO}(\mathcal{T}(M_n))$  is **non-unital**, **non-nuclear**.
- ▶ We have quotient maps

$$\mathcal{O}_{\mathcal{T}(M_n)} \rightarrow C(\mathbb{T})^{\star n} \quad (\text{non-unital free product})$$

$$S_{ij} \mapsto \delta_{ij} z_i.$$

$$\mathcal{O}_{\mathcal{T}(M_n)} \rightarrow U^{nc}(n)$$

$$S_{ij} \mapsto u_{ij}.$$

where  $U^{(nc)}(n) = C^* \left( u_{ij}, 1 \leq i, j \leq n \mid U = [u_{ij}] \text{ unitary} \right)$   
is **Brown's Universal noncommutative unitary algebra**.

- ▶ In fact  $M_n(\mathbb{FO}(\mathcal{T}(M_n))^+) \cong M_n \star (C(S^1) \oplus \mathbb{C})$ .

# Basic Examples, Basic Properties

**Example 2: Diagonal Quantum Graph**  $\mathcal{G} = \mathcal{D}(M_2)$

$$A_d(e_{ij}) = \delta_{i,1} 2e_{ij}.$$

$$\mathbb{FO}(\mathcal{D}(M_2)) = C^* \left( S_{ij}, 1 \leq i, j \leq 2 \mid \begin{array}{l} S = [S_{ij}] \text{ is a partial isometry} \\ \sum_l S_{li}^* S_{lj} = \delta_{\{i=j=1\}} \sum_l 2S_{1l} S_{1l}^* \end{array} \right)$$

- ▶ Get  $S_{l_2}^* S_{l_2} = 0 \implies S_{l_2} = 0$ .
- ▶ Thus, the canonical map  $S : B \rightarrow \mathcal{O}_{\mathcal{G}}$  need not be injective!

## Basic Examples, Basic Properties

**Example 3:** Complete Quantum Graph  $\mathcal{G} = \mathcal{K}(M_n)$ .

$$\mathbb{FO}(\mathcal{K}(M_n)) = C^* \left( S_{ij}, 1 \leq i, j \leq n \mid \begin{array}{l} S = [S_{ij}] \text{ is a partial isometry} \\ \sum_l S_{li}^* S_{lj} = \delta_{i,j} \left( n \sum_{x,l} S_{xl} S_{xl}^* \right) \end{array} \right)$$



## Basic Examples, Basic Properties

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- ▶ What is this algebra? Is it unital? Nuclear? ...
- ▶ We have a quotient map

$$\mathbb{FO}(\mathcal{K}(M_n)) \rightarrow \mathcal{O}_{n^2} = \mathcal{O}(K_{n^2}) \quad (\text{the Cuntz algebra})$$

$$S_{ij} \mapsto n^{-1/2} \hat{S}_{ij} \quad (\hat{S}_{ij} = \text{standard Cuntz isometry}).$$

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- ▶ **Very surprising/mysterious theorem:** The above homomorphism is actually an **isomorphism**  
 $\mathbb{FO}(\mathcal{K}(M_n)) \cong \mathcal{O}_{n^2}$ .
- ▶ So  $\mathbb{FO}(\mathcal{K}(M_n))$  is indeed unital, nuclear, etc.
- ▶ **The proof:** Uses the notion of a **quantum isomorphism** between quantum graphs (aka **superdense coding** in QIT):  
 $\mathcal{K}(M_n)$  and  $K_{n^2}$  turn out to be quantum isomorphic.

# Conclusion

Quantum graphs and Quantum Cuntz-Krieger Algebras seem to have some new and interesting properties! Lots of work to be done!

THANKS FOR LISTENING.