

# Lecture 21 - Jordan Algebras and Projective Spaces

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*References:*

*Jordan Operator Algebras. H. Hanche-Olsen and E. Stormer*

*The Octonions. J. Baez*

## 1 Jordan Algebras

### 1.1 Definition and examples

In the 1930's physicists, looking for a larger context in which to place quantum mechanics, settled on the following axioms for an algebra of observables:

- i)*  $\mathcal{A}$  is a vector space over  $\mathbb{R}$
- ii)* The product  $\circ$  is bilinear
- iii)* The product is commutative:  $A \circ B = B \circ A$
- iv)* The product obeys the *Jordan identity*:  $A \circ (B \circ (A \circ A)) = (A \circ B) \circ (A \circ A)$

An algebra obeying these axioms is called a *Jordan algebra*. A Jordan algebra is called *formally real* if it satisfies

- v)*  $A_1 \circ A_1 + \dots + A_k \circ A_k = 0$  if and only if  $A_1 = \dots = A_k = 0$ .

The motivating example was the set of physical observables; in the finite dimensional case, these are the hermitian  $n \times n$  matrices,  $\mathfrak{h}_n(\mathbb{C})$  with the product

$$A \circ B = \frac{1}{2}(AB + BA). \tag{1}$$

Two things can be noted here: this is formally real (unlike the full matrix algebra  $\mathbb{C}(n)$ ), and the Jordan product comes from the product in a larger associative algebra. A Jordan algebra whose product comes from an associative algebra via (1) is called a *special Jordan algebra*. The goal of the early explorers of this field was to come up with a broader theory of general Jordan algebras, into which  $\mathfrak{h}_n(\mathbb{C})$  might fit as some sort of special case.

Three series of finite dimensional formally real special Jordan algebras are

- The  $n \times n$  real symmetric matrices  $\mathfrak{h}_n(\mathbb{R})$
- The  $n \times n$  complex Hermitian matrices  $\mathfrak{h}_n(\mathbb{C})$
- The  $n \times n$  quaternionic Hermitian matrices  $\mathfrak{h}_n(\mathbb{H})$

In addition there is are the algebras of *spin factors*:

- $\mathbb{R}^n \oplus \mathbb{R}$  with product  $(X, a) \circ (Y, b) = (bX + aY, \langle X, Y \rangle + ab)$

## 1.2 Additional Facts

Normally we have left- and right- multiplication operators. Due to commutativity, we simple define the Jordan multiplication operator

$$T_a(b) = a \circ b. \quad (2)$$

The algebra of multiplication operators is associative (trivially); however due to non-associativity

$$\begin{aligned} T_a T_b &\neq T_{ab} \\ T_a T_b &\neq T_b T_a. \end{aligned} \quad (3)$$

Two elements  $a, b$  are said to *operator commute* if  $T_a T_b = T_b T_a$ . The *center* of a Jordan algebra is the set of all elements that operator commute. The Jordan identity can be expressed

$$T_a T_{a^2} = T_{a^2} T_a \quad (4)$$

or  $[T_a, T_{a^2}] = 0$ . The Jordan identity implies the following identity, called the *linearized Jordan identity*:

$$[T_a, T_{b \circ c}] + [T_b, T_{c \circ a}] + [T_c, T_{a \circ b}] = 0 \quad (5)$$

The first substantial theorem is that Jordan algebras are power-associative.

**Theorem 1.1** *Given a Jordan algebra  $\mathcal{J}$ , we have the following identities:*

$$\begin{aligned} a^m a^n &= a^{m+n} \\ T_{a^m} T_{a^n} &= T_{a^n} T_{a^m} \end{aligned} \quad (6)$$

□

We define a triple product

$$\{ABC\} = (A \circ B) \circ C + (B \circ C) \circ A - (A \circ C) \circ B. \quad (7)$$

In a special Jordan algebra one checks this is

$$\{ABC\} = \frac{1}{2}(ABC + CBA) \quad (8)$$

We define the operator  $U_{a,c}$  by

$$\begin{aligned} U_{a,c}(b) &\triangleq \{abc\} \\ U_{a,c} &= T_a T_c + T_c T_a - T_{a \circ c}. \end{aligned} \quad (9)$$

and the operator  $U_a = U_{a,a}$ , which is

$$U_{a,a} = 2(T_a)^2 - T_{a^2}. \quad (10)$$

**Theorem 1.2** *The Following identities hold in any Jordan algebra:*

$$\begin{aligned} (U_a)^n &= U_{a^n} \\ 2T_{a^t} U_{a^m, a^n} &= U_{a^{m+t}, a^n} + U_{a^m, a^{n+t}} \\ 2U_{a^m, a^n} T_{a^t} &= U_{a^{m+t}, a^n} + U_{a^m, a^{n+t}} \end{aligned} \quad (11)$$

A surprising fact is MacDonald's theorem

**Theorem 1.3 (MacDonald)** *Assume  $P(x, y, z)$  is any polynomial in three variables that has degree at most 1 in  $z$ . If  $P(x, y, z) = 0$  for all  $x, y, z$  in all special Jordan algebras,  $P(x, y, z) = 0$  in all Jordan algebras.*

**Theorem 1.4 (Shirshov-Cohn)** *Any Jordan algebra generated by two elements is special. In the unital case, any Jordan algebra generated by two elements and 1 is special.*

**Theorem 1.5 (Central Theorem)** *The center of any Jordan algebra is an associative algebra.*

**Theorem 1.6 (Complete Reducibility)** *If  $\mathcal{I} \subset \mathcal{J}$  is an ideal in a finite dimensional Jordan algebra, then  $\mathcal{J} = \mathcal{I} \oplus \mathcal{I}'$ .*

### 1.3 Projectors and the Classification

A projection in a Jordan algebra is an element  $p$  with  $p^2 = p$ . Lie algebras can be considered algebras of derivations; in an analogous sense, Jordan algebras can be considered algebras over projectors. As an example of this idea, consider  $\mathfrak{h}_2(\mathbb{C})$ . If  $(x, y) \in \mathbb{C}^2$  has  $|x|^2 + |y|^2 = 1$ , then

$$\begin{pmatrix} \bar{x} \\ \bar{y} \end{pmatrix} (x, y) = \begin{pmatrix} |x|^2 & \bar{x}y \\ \bar{y}x & |y|^2 \end{pmatrix} \quad (12)$$

is a projector in  $\mathfrak{h}_2(\mathbb{C})$ . Linear combinations of these constitute the entire algebra.

**Theorem 1.7** *The finite dimensional formally real Jordan algebras are the following:*

- $\mathfrak{h}_n(\mathbb{R})$ ,  $n \geq 3$
- $\mathfrak{h}_n(\mathbb{C})$ ,  $n \geq 3$
- $\mathfrak{h}_n(\mathbb{H})$ ,  $n \geq 3$
- the spin factors  $J_n$
- the Albert algebra, also known as the exceptional Jordan algebra.

The Albert algebra is  $\mathfrak{h}_3(\mathbb{O})$  with Jordan product  $a \circ b = \frac{1}{2}(ab + ba)$ .

*Pf.* See the reference. The idea is that Jordan algebras have so many projectors, that projections onto the entries of an  $n \times n$  matrix can be found.  $\square$

This theorem was a disappointment to the physicists of the 1950's. It meant there was no "general type" formally real Jordan algebra, only the special algebras and a single exceptional algebra.

### 1.4 Projectors and Projective Spaces

Given two projectors  $p : \mathbb{K} \rightarrow V_p$ ,  $q : \mathbb{K} \rightarrow V_q$  we say  $q \leq p$  if  $V_q \subseteq V_p$  (with strict inequality corresponding to strict inequality). Given  $p$ , we can generally construct chains of inequalities of projectors:

$$0 = p_0 < p_1 < \dots < p_n = p \quad (13)$$

Given  $p \in \mathcal{J}$ , the largest possible such  $n$  is called the *rank* of  $p$  in  $\mathcal{J}$ .

The projectors in  $\mathcal{J}$  can be used to construct projective spaces. The points are the rank-1 projectors, the lines are the rank-2 projectors, and so forth. Inclusion is the inequality, so

given a rank 1 projector  $p_1$  and a rank 2 projector  $p_2$ , we can say the point  $p_1$  is in the line  $p_2$  provided  $p_1 < p_2$ .

As one expects, this process carried out in  $\mathfrak{h}_n(\mathbb{R})$  gives  $\mathbb{R}P^{n-1}$ , in  $\mathfrak{h}_n(\mathbb{C})$  gives  $\mathbb{C}P^{n-1}$ , and in  $\mathfrak{h}_n(\mathbb{H})$  gives  $\mathbb{H}P^{n-1}$ . Carried out in the spin factors, one obtains a series of 1-dimensional projective planes. The algebra  $\mathfrak{h}_2(\mathbb{O})$  is a spin factor, and has associated to it the octonionic projective line, a manifold of 8 real dimensions diffeomorphic to  $\mathbb{S}^8$ . The algebra  $\mathfrak{h}_3(\mathbb{O})$  produces octonionic projective plane, a manifold of 16 real dimensions.

If  $\mathbb{P}$  is a projective space generated by a one of the  $\mathfrak{h}_n(\mathbb{K})$ , we have that any  $\vec{x} = (x_1, \dots, x_n)$  determines a projection by first normalizing to unit length, and then defining

$$p_{\vec{x}} = \frac{\vec{x} \vec{x}^T}{\vec{x}^T \vec{x}} \in \mathfrak{h}_n(\mathbb{K}) \quad (14)$$

gives a point in  $\mathbb{K}P^{n-1}$ .

## 2 Octonionic Projective Spaces

### 2.1 Geometric construction

Due to non-associativity, it is impossible to produce any kind of  $\mathbb{O}P^n$  for  $n \geq 3$ . However 1- and 2-dimensional octonionic projective planes have reasonable constructions.

By analogy with other octonionic lines, we can define  $\mathbb{O}P^1$  as the 1-point compactification of  $\mathbb{O} \approx \mathbb{R}^8$ , which is  $\mathbb{S}^8$ . The building of  $\mathbb{O}P^2$  can take place via the octonionic Hopf map. We have the unit sphere

$$\mathbb{S}^{15} \subset \mathbb{O}^2 \quad (15)$$

and therefore a map  $\mathbb{S}^{15} \rightarrow \mathbb{S}^8$  given by  $(x, y) \mapsto xy^{-1}$ . This gives  $\mathbb{S}^{15}$  as an  $\mathbb{S}^7$ -bundle over  $\mathbb{S}^8$ . We define

$$\mathbb{O}P^2 = (\mathbb{D}_1 \cup \mathbb{O}P^1) / \text{equiv} \quad (16)$$

where  $\mathbb{D}_1$  is the unit disk in  $\mathbb{O}^2$ , and the equivalence is the the Hopf map on  $\partial\mathbb{D}_1 = \mathbb{S}^{15} \rightarrow \mathbb{S}^8 \approx \mathbb{O}P^1$ .

### 2.2 The $\mathbb{O}P^n$ as projective spaces

In terms of the projective spaces  $\mathbb{K}P^n$  defined above, two vectors

$$\vec{x} = (x_1, \dots, x_{n+1}), \quad \vec{y} = (y_1, \dots, y_{n+1}) \quad (17)$$

determine the same point iff  $p_{\vec{x}} = p_{\vec{y}}$ . This determines an equivalence relation on these vectors:

$$\vec{x} \text{ equiv } \vec{y} \quad \text{iff} \quad p_{\vec{x}} = p_{\vec{y}}. \quad (18)$$

An equivalence class will be denoted

$$[\vec{x}] = [x_1, \dots, x_n] \quad (19)$$

If the base field is  $\mathbb{O}$ , then

$$\begin{aligned} [x_1, x_2] &\neq \{ (\alpha x_1, \alpha x_2) \mid \alpha \in \mathbb{O} \} \\ [x_1, x_2, x_3] &\neq \{ (\alpha x_1, \alpha x_2, \alpha x_3) \mid \alpha \in \mathbb{O} \} \end{aligned} \quad (20)$$

### 3 The Punchline

In the non-commutative case (let alone the non-associative case) it is difficult to define a determinant. However in the  $\mathfrak{h}_2(\mathbb{K})$  case it is always possible:

$$\det \begin{pmatrix} \alpha + \beta & x \\ \bar{x} & \alpha - \beta \end{pmatrix} = \alpha^2 - \beta^2 - |x|^2 \quad (21)$$

which gives a Lorentzian metric on  $\mathbb{R}^{n+1} \oplus \mathbb{R}$ , and an automorphism

$$SL(2, \mathbb{K}) \rightarrow SO_0(1, 3), \quad (22)$$

assuming, of course, that  $SL(2, \mathbb{H})$  and  $SL(2, \mathbb{O})$  can be independently defined! Indeed they can, and we have the expected isomorphisms

$$\begin{aligned} SL(2, \mathbb{R}) &\approx Spin(1, 2) \\ SL(2, \mathbb{C}) &\approx Spin(1, 3) \\ SL(2, \mathbb{H}) &\approx Spin(1, 5) \\ SL(2, \mathbb{O}) &\approx Spin(1, 9) \end{aligned} \quad (23)$$

The first two of these we saw in a previous lecture.

There will be more about this in the homeworks—in fact  $SL(n, \mathbb{H})$  and  $\mathfrak{sl}(n, \mathbb{H})$  always have a reasonable definition, and also  $SL(2, \mathbb{O})$ ,  $\mathfrak{sl}(2, \mathbb{O})$  can be defined.

In the case of  $A \in \mathfrak{h}_3(\mathbb{O})$  we have the following definition of the determinant:

$$\begin{aligned} \det A &= \frac{1}{3}Tr(A^3) - \frac{1}{2}Tr(A^2)Tr(A) + \frac{1}{6}(Tr(A))^3 \\ \det \begin{pmatrix} \alpha & \bar{z} & \bar{y} \\ z & \beta & x \\ y & \bar{x} & \gamma \end{pmatrix} &= \alpha\beta\gamma - (\alpha|x|^2 + \beta|y|^2 + \gamma|z|^2) + 2Re(xyz). \end{aligned} \quad (24)$$

The automorphisms that preserve the Hermitian inner product on  $\mathbb{O}^2$  and  $\mathbb{O}^3$  are the groups  $SU(2, \mathbb{O})$  and  $SU(3, \mathbb{O})$  corresponding to the groups that preserve the Jordan product in these dimensions. We have

$$SU(2, \mathbb{O}) \approx SO(9), \tag{25}$$

the isometry group of  $\mathbb{O}P^1 \approx \mathbb{S}^8$ .

The group of automorphisms of  $\mathfrak{h}_3(\mathbb{O})$  that preserve the Jordan product is a real form of  $F_4$ . In an appropriate sense, we may therefore regard

$$F_4 \approx SU(3, \mathbb{O}). \tag{26}$$

The group of maps  $\mathfrak{h}_3(\mathbb{O}) \rightarrow \mathfrak{h}_3(\mathbb{O})$  that preserve the determinant is a real form of  $E_6$ . In an appropriate sense, we may therefore regard

$$E_6 \approx SL(3, \mathbb{O}). \tag{27}$$

We will give justifications next time, and begin the constructions that culminate in the “magic square.”