

# Lecture 18 - Clifford Algebras and Spin groups

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*Reference: Lawson and Michelsohn, Spin Geometry.*

## 1 Universal Property

If  $V$  is a vector space over  $\mathbb{R}$  or  $\mathbb{C}$ , let  $q$  be any quadratic form, meaning a map  $q : V \rightarrow W$  with  $q(\alpha v) = |\alpha|^2 q(v)$  for scalars  $\alpha$ . Given a pair  $(V, q)$  we can construct the Clifford algebra  $Cl(V, q)$  as the quotient of the tensor algebra  $\otimes V$  by the ideal generated by  $v \otimes v + q(v)$ , as  $v$  ranges over  $V$ .

Note that if  $q = 0$ , then  $Cl(V, q) = \wedge V$ .

Ordinarily we consider the case that  $q$  is nondegenerate. Over the complex numbers, nondegenerate quadratic forms are all equivalent up to change of basis, so the only invariant in the dimension of the space  $V$ . Setting  $\dim_{\mathbb{C}}(V) = n$  we obtain the algebras  $Cl_n$ . In the real setting, besides the dimension of the vector space  $\dim_{\mathbb{R}}(V) = n$ , the only other invariant of a quadratic form is its signature, which is characterized by a pair of non-negative integers  $(r, s)$  where  $r + s = n$ . We obtain the Clifford algebras  $Cl_{r,s}$ .

Clifford algebras have the following universal property. If  $(V, q)$  is a vector space (over  $\mathbb{R}$  or  $\mathbb{C}$ ) with quadratic form  $q$  and  $\varphi : V \rightarrow GL(W, W)$  is a linear map so that  $\varphi(v)\varphi(v) = -q(v)^2$ , then  $\varphi$  extends uniquely to a Clifford representation map  $\varphi : Cl(V, q) \rightarrow GL(W, W)$ .

## 2 Reason for Studying Clifford Algebras

Let  $\mathbb{D}$  be a normed division algebra over reals, with norm  $\|\cdot\|$ . We obtain an inner product

$$\langle x, y \rangle = \frac{1}{2} (\|x + y\|^2 - \|x\|^2 - \|y\|^2). \quad (1)$$

Since  $\|ax\| = \|a\|\|x\|$ , we have that  $L_a$  is an orthogonal transformation whenever  $\|a\| = 1$ .

Consider the subspace of vectors in  $\mathbb{D}$  orthogonal to 1; we shall call this  $Im(\mathbb{D})$ . The unit sphere in  $\mathbb{D}$  is a group of orthogonal transformation of  $\mathbb{D}$ , and the tangent space to the unit sphere at 1, its Lie algebra, can be identified with  $Im(\mathbb{D})$ . Therefore  $L_a$  for  $a \in Im(\mathbb{D})$  is skew-symmetric. If  $a$  is also a unit vector, then

$$\langle 1, a^2 \rangle = -\langle a, a \rangle = -1 \quad (2)$$

so that  $a^2 = -1$ . If  $dim(\mathbb{D}) = n$ , this gives a faithful representation of  $Cl_{n-1}$  on  $\mathbb{D}$ .

### 3 Periodicity

**Proposition 3.1** *We have the following algebra isomorphisms*

$$\begin{aligned} Cl_{n,0} \otimes Cl_{0,2} &\approx Cl_{0,n+2} \\ Cl_{0,n} \otimes Cl_{2,0} &\approx Cl_{n+2,0} \\ Cl_{r,s} \otimes Cl_{1,1} &\approx Cl_{r+1,s+1} \end{aligned} \quad (3)$$

*Pf.*

We do only the first isomorphism; the rest are entirely analogous. Consider the vector space  $\mathbb{R}^{n+2}$  with orthonormal basis  $e_1, \dots, e_{n+2}$ , with negative definite inner product. Also let  $f_1, \dots, f_n$  and  $g_{n+1}, g_{n+2}$  be orthonormal bases of  $\mathbb{R}^n$  with positive definite inner product, and  $\mathbb{R}^2$  with negative definite inner product. Let  $F : \mathbb{R}^{n+2} \rightarrow Cl_{n,0} \otimes Cl_{0,2}$  be the map

$$F(e_i) = \begin{cases} f_i \otimes g_{n+1}g_{n+2} & , \quad \text{if } i \in \{1, \dots, n\} \\ 0 \otimes g_i & , \quad \text{if } i \in \{n+1, n+2\} \end{cases} \quad (4)$$

Noting that  $F(e_i)F(e_j) = -F(e_j)F(e_i)$  when  $i \neq j$  and  $F(e_i)^2 = 1$ , this extends to the desired algebra isomorphism.  $\square$

**Theorem 3.2 (“Bott periodicity”)** *We have the following isomorphisms*

$$\begin{aligned} Cl_{n+8,0} &\approx Cl_{n,0} \otimes_{\mathbb{R}} Cl_{8,0} \approx Cl_{n,0} \otimes_{\mathbb{R}} \mathbb{R}(16) \\ Cl_{0,n+8} &\approx Cl_{0,n} \otimes_{\mathbb{R}} Cl_{0,8} \approx Cl_{0,n} \otimes_{\mathbb{R}} \mathbb{R}(16) \\ Cl_{n+2} &\approx Cl_n \otimes_{\mathbb{C}} Cl_2 \approx Cl_n \otimes_{\mathbb{C}} \mathbb{C}(2) \end{aligned} \quad (5)$$

*Pf.*

Noting that  $\mathbb{H}^{\otimes 2} \approx \mathbb{R}(4)$ , we can use the previous lemma to obtain

$$\begin{aligned} Cl_{n+8,0} &\approx Cl_{n,0} \otimes (Cl_{0,2} \otimes Cl_{2,0})^{\otimes 2} \\ &\approx Cl_{n,0} \otimes \mathbb{H} \otimes \mathbb{H} \otimes \mathbb{R}(2) \otimes \mathbb{R}(2) \\ &\approx Cl_{n,0} \otimes \mathbb{R}(4) \otimes \mathbb{R}(4) \\ &\approx Cl_{n,0} \otimes \mathbb{R}(16) \end{aligned} \quad (6)$$

The other cases are similar.  $\square$

## 4 Pin, Spin, and Orthogonal Groups

The “group of units” in  $Cl_{p,q}$ , denoted  $Cl_{p,q}^\times$ , is the group of all invertible elements. The associated Lie algebra is denoted  $\mathfrak{cl}_{p,q}^\times$ . In fact we have  $\mathfrak{cl}_n^\times \approx Cl_{p,q}$  with the Lie bracket  $[x, y] = xy - yx$ .

We have an adjoint map  $Ad : Cl_{p,q}^\times \rightarrow Aut(Cl_{p,q})$

$$Ad_\varphi(x) = \varphi x \varphi^{-1} \quad (7)$$

whose derivative (in the appropriate sense) is the adjoint map  $ad : \mathfrak{cl}_{p,q}^\times \rightarrow Der(Cl_{p,q})$

$$ad_y(x) = [y, x]. \quad (8)$$

Let  $(V, q)$  be an  $n$ -dimensional vector space with non-degenerate quadratic form  $q$ . Polarizing the quadratic form  $q$  of signature  $(p, q)$  we obtain an inner product  $(\cdot, \cdot)$ . Given a vector  $v \in V$ , let  $\sigma_v$  be the reflection about the hyperplane orthogonal to  $v$ , so that

$$\sigma_v(x) = x - 2 \frac{(x, v)}{(v, v)} v. \quad (9)$$

An important subgroup of  $Cl^\times(V, q)$  is the group  $P(V, q)$  generated by elements  $v \in V$  with  $q(v) \neq 0$ . Quotienting out by constants, we obtain the Pin group. Specifically,  $Pin(V, q)$  (or  $Pin_{p,q}$ ) is the group generated by elements  $v \in V$  with  $q(v) = \pm 1$ . Further define the spin groups to be

$$Spin(V, q) = Pin(V, q) \cap Cl^0(V, q). \quad (10)$$

**Proposition 4.1** *The adjoint representation of  $Pin(V, q)$  on  $Cl(V, q)$  restricts to a representation of  $Pin(V, q)$  on  $V$ . For any  $v \in V \cap Cl^\times(V, q)$  we have*

$$-Ad_v|_V = \sigma_v. \quad (11)$$

Thus we have found a map

$$Pin_{p,q} \longrightarrow O(p, q). \quad (12)$$

As subsets of the Clifford algebras, the spin groups contain only monomials with even numbers of generators. Thus we have found a map

$$Spin_{p,q} \longrightarrow SO(p, q). \quad (13)$$

The adjoint representation is somewhat inadequate, due to the negative. For instance, if  $V$  is odd dimensional, then  $Ad_v$  always preserves orientation. We introduce the twisted adjoint map

$$\widetilde{Ad}_\varphi(x) = \varphi x \alpha(\varphi^{-1}) \quad (14)$$

where  $\alpha$  is the parity involution. We have that  $\widetilde{Ad}_v$  is a representation (as  $\widetilde{Ad}_\varphi \widetilde{Ad}_\eta = \widetilde{Ad}_{\varphi\eta}$ ) and that

$$\widetilde{Ad}_v = \sigma_v. \quad (15)$$

when  $v \in V$ .

**Theorem 4.2** *The  $\widetilde{Ad}$  maps*

$$\begin{aligned} Pin_{r,s} &\longrightarrow O(r,s) \\ Spin_{r,s} &\longrightarrow SO(r,s) \end{aligned} \quad (16)$$

are surjective. Their kernels are  $\mathbb{Z}_2 \approx \{\pm 1\}$ . Except in the cases  $(r,s) = (1,1), (1,0), (0,1)$ , these two-sheeted coverings are connected over any connected component of  $O(r,s)$  (resp.  $SO(r,s)$ ).

*Pf.* The only thing to prove is that 1 and  $-1$  can be joined by a path in  $Spin_{r,s} \subset Pin_{r,s}$ . Choosing orthogonal unit vectors  $e_1, e_2$  with  $q(e_1) = q(e_2)$  (either both  $+1$  or both  $-1$ ), we have that  $e_1 \cos(t) + e_2 \sin(t)$  and  $-e_1 \cos(t) + e_2 \sin(t)$  are unit vectors and

$$\begin{aligned} \gamma(t) &= (e_1 \cos(t) + e_2 \sin(t))(-e_1 \cos(t) + e_2 \sin(t)) \\ &= \pm (\cos^2(t) - \sin^2(t)) + 2e_1 e_2 \cos(t) \sin(t) \\ &= \pm \cos(2t) + e_1 e_2 \sin(2t) \end{aligned} \quad (17)$$

is a path from 1 to  $-1$  in  $Spin_{r,s}$ . □

## 5 Spin Representations

### 5.1 Complete Reducibility

A basic theme in the study of Clifford representation is the reduction of the representation to a Pin or Spin representation. For instance, any representations of a Clifford algebra is completely reducible. To see this, note that any Clifford representation  $V$  is a Pin representation, so  $V$  is completely reducible as a Pin representation. But  $Pin_{r,s}$  contains a basis of  $Cl_{r,s}$ , so the action of the Clifford algebra preserves the factors.

## 5.2 Fundamental Irreducible Clifford Representations

We have the following table

$n$	$Cl_n$	$Cl_n^0$
0	$\mathbb{R}$	–
1	$\mathbb{C}$	$\mathbb{R}$
2	$\mathbb{H}$	$\mathbb{C}$
3	$\mathbb{H} \oplus \mathbb{H}$	$\mathbb{H}$
4	$\mathbb{H}(2)$	$\mathbb{H} \oplus \mathbb{H}$
5	$\mathbb{C}(4)$	$\mathbb{H}(2)$
6	$\mathbb{R}(8)$	$\mathbb{C}(4)$
7	$\mathbb{R}(8) \oplus \mathbb{R}(8)$	$\mathbb{R}(8)$
8	$\mathbb{R}(16)$	$\mathbb{R}(8) \oplus \mathbb{R}(8)$

(18)

This table can be used to give the irreducible representations of  $Cl_n$ . In particular, with  $n \not\equiv 3 \pmod{4}$  we have a single irreducible representation, and with  $n \equiv 3 \pmod{4}$ , we have two representations. For  $n \equiv 3 \pmod{8}$ , these representations are inequivalent.

However, we proceed more theoretically. Consider the Clifford group, the finite group  $E_n$  generated by an orthonormal basis  $e_1, \dots, e_n$  of  $\mathbb{R}^n$ . We have that the group algebra  $\mathbb{R}E_n$  is nearly the Clifford algebra: in particular

$$Cl_n \approx \mathbb{R}E_n / \{(-1) + 1 = 0\}. \quad (19)$$

Note that

$$E_n = \left\{ \pm \prod_{i=1}^m e_i^{k_i} \mid k_i = \pm 1 \right\} \quad (20)$$

is a subgroup of  $Cl_n$  with  $2^{n+1}$  elements. Set  $E_n^0 = E_n \cap Cl_n^0$ , so that  $E_n^0$  has  $2^n$  elements. Finally, if  $n$  is even, then the center of  $E_n$  is  $\{\pm 1\}$  and if  $n$  is odd, its center is  $\{\pm 1, \pm \omega\}$  where  $\omega = e_1 \dots e_n$ .

Now consider that each element of  $E_n$  acts on a 1-dimensional subspace, so we have a  $2^{n+1}$ -dimensional representation of  $\mathbb{R}E_n$ , which passes to a  $2^n$  dimensional representation of  $Cl_n$ , which we shall call  $W_n$ .

**Proposition 5.1 (Clifford Representations)** *In the cases  $n \equiv 0, 1, 2 \pmod{4}$ , the Clifford algebra  $Cl_n$  has a single irreducible representation  $W_n$  of dimension  $2^n$ . In the case that  $n \equiv 3 \pmod{4}$ , the Clifford algebra  $Cl_n$  has two inequivalent irreducible representations  $W_n^+$  and  $W_n^-$  of dimensions  $2^{n-1}$ .*

*Pf.* We consider the case of  $n \equiv 3 \pmod{4}$ . Let  $W_n$  be the algebra constructed above. We know that  $\omega$  is central and  $\omega^2 = 1$ . Therefore  $\pi^\pm = \frac{1}{2}(1 \pm \omega)$  splits the representation:

$W_n^\pm = \pi^\pm W_n$ . To see that these representations are inequivalent, notice that  $\omega$  acts as 1 on  $W_n^+$  and  $-1$  on  $W_n^-$ .  $\square$

Notice that if  $n = 4m$ , then  $\omega^2 = 1$  and  $W_n$  also splits as a vector space:  $W_n = W_n^+ \oplus W_n^-$  where  $W_n^\pm = \pi^\pm W$ . However  $\pi^\pm$  is not central, so the factors are exchanged by elements of  $Cl_n^1$ . But  $Cl_n^0$  preserves the factors.

We can restrict these representations  $Cl_n$  on  $W_n$  to  $Spin_n \subset Cl_n^0 \subset Cl_n$ ; we denote the resulting representation  $S_n$ .

**Theorem 5.2 (Fundamental Spin Representations)** *If  $n \equiv 1, 2, 3 \pmod{4}$  then  $S_n$  is the direct sum of two identical representations, which we call  $\Delta_n$ . If  $n \equiv 0 \pmod{4}$  then  $S_n$  is the direct sum of two different, non-isomorphic representations; we denote  $S_n = \Delta_n = \Delta_n^+ \oplus \Delta_n^-$  where  $\Delta_n^\pm = \pi_{n-1}^\pm S_n$ .*

*Pf.* Note that the parity involution  $\alpha$  preserves  $Spin_n \subset Cl_n^0 \subset Cl_n$ , and that the projectors  $\pi^0 = \frac{1}{2}(1 + \alpha)$ ,  $\pi^1 = \frac{1}{2}(1 - \alpha)$  commute and sum to unity. Thus the representation of  $Spin_n$  on  $Cl_n$  splits into representations on  $Cl_n^0$  and  $Cl_n^1$ .

Case  $n = 2m + 1$  is odd. We have that  $\omega \notin spin(n)$ . If  $m$  is even ( $n \equiv 1 \pmod{4}$ ) then the representation on  $Cl_n^0$  is irreducible because  $Cl_n^0 \approx Cl_{n-1}$  is irreducible. Further, the factors  $Cl_n^0$  and  $Cl_n^1$  are exchanged by  $\omega$ , which is central; hence the representations on  $Cl_n^0$  and  $Cl_n^1$  are equivalent. If  $m$  is odd ( $n \equiv 3 \pmod{4}$ ) then we have a secondary decomposition  $Cl_n = Cl_n^+ \oplus Cl_n^-$ , where  $Cl_n^0$  sits in the sum diagonally. However  $Cl_n^0 \approx Cl_{n-1}$  is irreducible, and projects onto either factor isomorphically. Thus again we obtain two equivalent representations.

Case  $n = 2m$  is even. We have that  $\omega \in spin(n)$ , and that  $\omega$  is central in  $Spin(n)$  (although it is not central in  $Cl_n$ ). If  $m$  is odd so  $n \equiv 2 \pmod{4}$  then we get the usual even-odd splitting.

If  $m$  is even so  $n \equiv 0 \pmod{4}$  then  $\omega^2 = 1$  and we can decompose  $Cl_n$  into two factors, corresponding to the  $\pm 1$  eigenvalues on  $\omega$ . These factors are clearly inequivalent representations, and  $\omega \in Spin(n)$  acts differently. Also,

$\square$

An element of a fundamental spin representation  $v \in \Delta_n$  is called a spinor.