



MATH 241 AP EXAM

SPRING 2013

NAME:

Please *turn off and put away all electronic devices*. You may use both sides of a 8.5" \times 11" sheet of paper for handwritten notes while you take this exam. No calculators, no course notes, no books, no help from your neighbors. **Show all work**. Please **clearly mark** your final answer. Remember to put your name at the top of this page. Good luck.

My signature below certifies that I have complied with the University of Pennsylvania's Code of Academic Integrity in completing this examination.

Your signature

QUESTION NUMBER	POINTS POSSIBLE	YOUR SCORE
1	10	
2	10	
3	10	
4	10	
5	10	

QUESTION NUMBER	POINTS POSSIBLE	YOUR SCORE
6	10	
7	10	
8	10	
9	10	
10	10	
TOTAL	100	

A PARTIAL TABLE OF INTEGRALS

$$\int_0^x u \cos nu \, du = \frac{\cos nx + nx \sin nx - 1}{n^2} \quad \text{for any real } n \neq 0$$

$$\int_0^x u \sin nu \, du = \frac{\sin nx - nx \cos nx}{n^2} \quad \text{for any real } n \neq 0$$

$$\int_0^x e^{mu} \cos nu \, du = \frac{e^{mx}(m \cos nx + n \sin nx) - m}{m^2 + n^2} \quad \text{for any real } n, m$$

$$\int_0^x e^{mu} \sin nu \, du = \frac{e^{mx}(-n \cos nx + m \sin nx) + n}{m^2 + n^2} \quad \text{for any real } n, m$$

$$\int_0^x \sin nu \cos mu \, du = \frac{m \sin nx \sin mx + n \cos nx \cos mx - n}{m^2 - n^2} \quad \text{for any real numbers } m \neq n$$

$$\int_0^x \cos nu \cos mu \, du = \frac{m \cos nx \sin mx - n \sin nx \cos mx}{m^2 - n^2} \quad \text{for any real numbers } m \neq n$$

$$\int_0^x \sin nu \sin mu \, du = \frac{n \cos nx \sin mx - m \sin nx \cos mx}{m^2 - n^2} \quad \text{for any real numbers } m \neq n$$

FORMULAS INVOLVING BESSEL FUNCTIONS

- Bessel's equation: $r^2 R'' + rR' + (\alpha^2 r^2 - n^2)R = 0$ – The only solutions of this which are bounded at $r = 0$ are $R(r) = cJ_n(\alpha r)$.

$$J_n(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!(k+n)!} \left(\frac{x}{2}\right)^{n+2k}.$$

$J_0(0) = 1$, $J_n(0) = 0$ if $n > 0$. z_{nm} is the m th positive zero of $J_n(x)$.

- Orthogonality relations:

$$\text{If } m \neq k \text{ then } \int_0^1 x J_n(z_{nm}x) J_n(z_{nk}x) \, dx = 0 \quad \text{and} \quad \int_0^1 x (J_n(z_{nm}x))^2 \, dx = \frac{1}{2} J_{n+1}(z_{nm})^2.$$

- Recursion and differentiation formulas:

$$\frac{d}{dx}(x^n J_n(x)) = x^n J_{n-1}(x) \quad \text{or} \quad \int x^n J_{n-1}(x) \, dx = x^n J_n(x) + C \quad \text{for } n \geq 1 \quad (1)$$

$$\frac{d}{dx}(x^{-n} J_n(x)) = -x^{-n} J_{n+1}(x) \quad \text{for } n \geq 0 \quad (2)$$

$$J'_n(x) + \frac{n}{x} J_n(x) = J_{n-1}(x) \quad (3)$$

$$J'_n(x) - \frac{n}{x} J_n(x) = -J_{n+1}(x) \quad (4)$$

$$2J'_n(x) = J_{n-1}(x) - J_{n+1}(x) \quad (5)$$

$$\frac{2n}{x} J_n(x) = J_{n-1}(x) + J_{n+1}(x) \quad (6)$$

- Modified Bessel's equation: $r^2 R'' + rR' - (\alpha^2 r^2 + n^2)R = 0$ – The only solutions of this which are bounded at $r = 0$ are $R(r) = cI_n(\alpha r)$.

$$I_n(x) = i^{-n} J_n(ix) = \sum_{k=0}^{\infty} \frac{1}{k!(k+n)!} \left(\frac{x}{2}\right)^{n+2k}.$$

FORMULAS INVOLVING ASSOCIATED LEGENDRE AND SPHERICAL BESSEL FUNCTIONS

- Associated Legendre Functions: $\frac{d}{d\phi} \left(\sin \phi \frac{dg}{d\phi} \right) + \left(\mu - \frac{m^2}{\sin^2 \phi} \right) g = 0$. Using the substitution $x = \cos \phi$, this equation becomes $\frac{d}{dx} \left((1-x^2) \frac{dg}{dx} \right) + \left(\mu - \frac{m^2}{1-x^2} \right) g = 0$. This equation has bounded solutions only when $\mu = n(n+1)$ and $0 \leq m \leq n$. The solution $P_n^m(x)$ is called an associated Legendre function of the first kind.

- Associated Legendre Function Identities:

$$P_n^0(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} (x^2 - 1)^n \text{ and } P_n^m(x) = (-1)^m (1-x^2)^{m/2} \frac{d^m}{dx^m} P_n(x) \text{ when } 1 \leq m \leq n$$

- Orthogonality of Associated Legendre Functions: If n and k are both greater than or equal to m ,

$$\text{If } n \neq k \text{ then } \int_{-1}^1 P_n^m(x) P_k^m(x) dx = 0 \text{ and } \int_{-1}^1 (P_n^m(x))^2 dx = \frac{2(n+m)!}{(2n+1)(n-m)!}.$$

- Spherical Bessel Functions: $(\rho^2 f')' + (\alpha^2 \rho^2 - n(n+1))f = 0$. If we define the spherical Bessel function $j_n(\rho) = \rho^{-\frac{1}{2}} J_{n+\frac{1}{2}}(\rho)$, then only solution of this ODE bounded at $\rho = 0$ is $j_n(\alpha\rho)$.

- Spherical Bessel Function Identity:

$$j_n(x) = x^2 \left(-\frac{1}{x} \frac{d}{dx} \right)^n \left(\frac{\sin x}{x} \right).$$

- Spherical Bessel Function Orthogonality: Let z_{nm} be the m -th positive zero of j_n .

$$\text{If } m \neq k \text{ then } \int_0^1 x^2 j_n(z_{nm}x) j_n(z_{nk}x) dx = 0 \text{ and } \int_0^1 x^2 (j_n(z_{nm}x))^2 dx = \frac{1}{2} (j_{n+1}(z_{nm}))^2.$$

ONE-DIMENSIONAL FOURIER TRANSFORM

$$\mathcal{F}[u](\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} u(x) e^{i\omega x} dx, \quad \mathcal{F}^{-1}[U](x) = \int_{-\infty}^{\infty} U(\omega) e^{-i\omega x} d\omega$$

TABLE OF FOURIER TRANSFORM PAIRS
 FOURIER TRANSFORM PAIRS FOURIER TRANSFORM PAIRS
 ($\alpha > 0$) ($\beta > 0$)

$u(x) = \mathcal{F}^{-1}[U]$	$U(\omega) = \mathcal{F}[u]$	$u(x) = \mathcal{F}^{-1}[U]$	$U(\omega) = \mathcal{F}[u]$
$e^{-\alpha x^2}$	$\frac{1}{\sqrt{4\pi\alpha}} e^{-\frac{\omega^2}{4\alpha}}$	$\sqrt{\frac{\pi}{\beta}} e^{-\frac{x^2}{4\beta}}$	$e^{-\beta\omega^2}$
$e^{-\alpha x }$	$\frac{1}{2\pi} \frac{2\alpha}{x^2 + \alpha^2}$	$\frac{2\beta}{x^2 + \beta^2}$	$e^{-\beta \omega }$
$u(x) = \begin{cases} 0 & x > \alpha \\ 1 & x < \alpha \end{cases}$	$\frac{1}{\pi} \frac{\sin \alpha\omega}{\omega}$	$2 \frac{\sin \beta x}{x}$	$U(\omega) = \begin{cases} 0 & \omega > \beta \\ 1 & \omega < \beta \end{cases}$
$\delta(x - x_0)$	$\frac{1}{2\pi} e^{i\omega x_0}$	$e^{-i\omega_0 x}$	$\delta(\omega - \omega_0)$
$\frac{\partial u}{\partial t}$	$\frac{\partial U}{\partial t}$	$\frac{\partial^2 u}{\partial t^2}$	$\frac{\partial^2 U}{\partial t^2}$
$\frac{\partial u}{\partial x}$	$-i\omega U$	$\frac{\partial^2 u}{\partial x^2}$	$(-i\omega)^2 U$
xu	$-i \frac{\partial U}{\partial \omega}$	$x^2 u$	$(-i)^2 \frac{\partial^2 U}{\partial \omega^2}$
$u(x - x_0)$	$e^{i\omega x_0} U$	$\frac{1}{2\pi} \int_{-\infty}^{\infty} f(s)g(x-s)ds$	FG

1. Beginning with the principle of conservation of energy and Fourier's law of heat conduction, derive the PDE that governs heat flow on a long thin rod assuming that density is a function of x but all other quantities (cross-sectional area, thermal conductivity, specific heat) are constant.

2. Solve the Laplace equation on the half disk $0 \leq \theta \leq \pi$, $0 \leq r \leq 1$, subject to the boundary conditions $\frac{\partial u}{\partial r}(1, \theta) = \theta + 2$ and $u(r, 0) = u(r, \pi) = 0$.

3. (a) Write the function $\frac{1}{2} - x$ as a Fourier sine series on the interval $[0, \frac{1}{2}]$.
- (b) Carefully plot the sum of the Fourier series on the interval $[-2, 2]$.
- (c) Use your answer from (a) to compute the Fourier cosine series of $x(1 - x)$ on the interval $[0, \frac{1}{2}]$.

4. Consider the displacement u of a string with source term $Q = -\alpha \frac{\partial u}{\partial t}$ (where α is a small, positive constant), given by

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2} - \alpha \frac{\partial u}{\partial t}.$$

- (a) What is the physical effect of the source term Q on the string?
- (b) Using separation of variables, solve the equation subject to the constraints

$$u(0, t) = 0, \quad u(L, t) = 0, \quad u(x, 0) = 0, \quad \frac{\partial u}{\partial t}(x, 0) = f(x).$$

- (c) In the case of part (b), what are the (circular) frequencies of vibration of the string?

5. Express the eigenvalue problem

$$e^{x^2} \frac{d^2 \phi}{dx^2} + x \frac{d\phi}{dx} = -\lambda x^2 \phi \text{ on } [1, 2] \text{ with } \phi(1) = \phi(2) = 0$$

in the standard Sturm-Liouville form. Explicitly verify that the Sturm-Liouville problem is regular, give the orthogonality relationship satisfied by the eigenfunctions, an asymptotic formula for the eigenvalues, and an approximate expression for the eigenfunctions themselves.

6. Write a series for the solution u of the modified heat equation

$$\frac{\partial u}{\partial t} + tu = k\nabla^2 u$$

on the unit ball in three dimensions (and all $t \geq 0$) subject to homogeneous Dirichlet boundary conditions. You need not relate the coefficients of the series to the initial data.

7. Let $u(x, t)$ be the solution of the initial/boundary value problem for the heat equation:

$$u_t = u_{xx} \quad \text{for } 0 < x < 1, \quad t > 0$$
$$u(x, 0) = 0, \quad \frac{\partial u}{\partial x}(0, t) = 1, \quad u(1, t) = 1.$$

(a) What is $\lim_{t \rightarrow \infty} u(x, t)$?

(b) Find $u(x, t)$.

(c) Use the equilibrium solution and one term from the Fourier series of $u(x, t)$ to estimate how long from $t = 0$ it will take for $u(\frac{1}{2}, t)$ to attain half its equilibrium value (in other words, estimate the value of T such that

$$u\left(\frac{1}{2}, t\right) < \frac{1}{2} \lim_{t \rightarrow \infty} u\left(\frac{1}{2}, t\right) \quad \text{for } t < T$$

and

$$u\left(\frac{1}{2}, t\right) > \frac{1}{2} \lim_{t \rightarrow \infty} u\left(\frac{1}{2}, t\right) \quad \text{for } t > T.$$

(Bonus points for explaining why this is a really good estimate.)

8. Consider the following wave equation on the interval $[0, L]$:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2},$$

$$u(x, 0) = \frac{\partial u}{\partial t}(x, 0) = 0, \quad \frac{\partial u}{\partial x}(0, t) = 0, \quad u(L, t) = \sin \alpha t$$

For what values of α is resonance possible? How does the situation change if the boundary condition at $x = 0$ is changed to $u(x, 0) = 0$? Note that you do not need to write down the solution of this PDE to answer this question.

9. Let

$$f(x) = \begin{cases} 1 & \text{for } 0 \leq x \leq 2 \\ 0 & \text{for } |x - 1| > 1 \end{cases}.$$

- (a) Compute the Fourier transform of f .
- (b) Solve initial-value problem for the wave equation

$$u_{tt} = \frac{1}{9}u_{xx}, \quad u(x, 0) = f(x), \quad u_t(x, 0) = 0$$

with $f(x)$ as given at the beginning of this problem.

- (c) Draw a reasonably accurate graph of $u(x, 3)$.

10. Solve the initial-value problem

$$u_t = 2u_{xx} + u_x, \quad u(x, 0) = e^{-x^2}$$

for $u(x, t)$, $t > 0$ and $-\infty < x < \infty$. Fully simplify your answer.

