

Singularity structure and massless dyons of pure $\mathcal{N} = 2, d = 4$ theories with A_r and C_r

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Disclaimer: all figures are mouse-drawn, not to scale.

Based on works w/ Keshav Dasgupta

Seiberg-Witten and Argyres-Douglas theories

- ▶ Seiberg-Witten curves, massless dyons (low ranks), and some aspects of singularity in moduli space were studied for
 - ▶ SU(2) w/o and w/ matter [Seiberg Witten 94]
 - ▶ SU(n) [Klemm Lerche Yankielowicz Theisen 94, Argyres Faraggi 94], SU(n) w/o and w/ matter [Hanany Oz 95]
 - ▶ SO(2r) [Brandhuber-Landsteiner 95] SO(2r+1) [Danielsson Sundborg 95], SO with matter [Hanany 95]
- ▶ The hyper-elliptic curve $y^2 = f(x)$ degenerates into a cusp form

$$y^2 = (x - a)^m \times \dots$$

when $m \geq 3$ branch points collide on x -plane. For SU(n) SW curves, this give exotic theories, discovered and studied [Argyres Douglas 95], where we have *mutually non-local* massless dyons.

(Under symplectic transformation, you cannot bring everything purely electronic.)

Main idea

Want to build Gaiotto's $\mathcal{N} = 2$ theory (see Yuji's talk) in F theory

- ▶ $\mathcal{N} = 2$ $Sp(2r)$ SW theories via r D3-branes
- ▶ compute massless dyon charges for pure $SU(r+1)$, $Sp(2r)$
- ▶ study wall crossing
- ▶ provide plethora of Argyres-Douglas (AD) theories

physics	\leftrightarrow	geometry
massless d.o.f.	\leftrightarrow	singularity $\Delta_x f = 0$
exotic theories (i. e. AD)	\leftrightarrow	$\Delta_u \Delta_x f = 0$

Double discriminant captures higher singularity of hyper-elliptic curves.

Discriminant

Δ_x : discriminant *with respect to x*

$$f_n(x) = \sum_{i=1}^n a_i x^i = a_n \prod_{i=1}^n (x - e_i)$$

$$\Delta_x(f_n(x)) = a_n^{2n-2} \prod_{i < j} (e_i - e_j)^2$$

Order of vanishing (multiplicity of roots) $\geq 2 \Leftrightarrow \Delta_x = 0$

Recipe for locating Argyres-Douglas loci

- ▶ Start with hyper-elliptic Seiberg-Witten curve

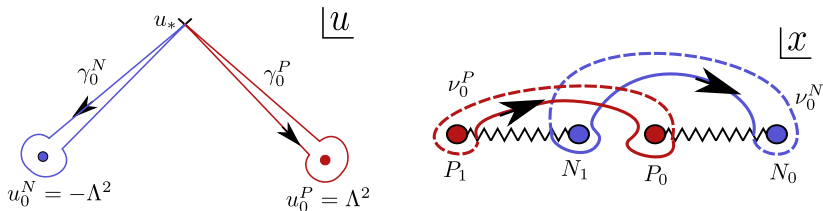
$$y^2 = f(x; u, v, \dots)$$
- ▶ Demanding $\Delta_x f = 0$ and $\Delta_u \Delta_x f = 0$ gives two massless BPS dyons. [Argyres Plesser Seiberg Witten '95]
- ▶ Check **order of vanishing** (o.o.v.) of each solution to $\Delta_u \Delta_x f = 0$
- ▶ If o.o.v. ≥ 3 , Argyres-Douglas loci: The hyperelliptic curve degenerates into a cusp-like singularity $y^2 = (x - a)^3 \times \dots$ and two mutually non-local dyons become massless. (checked up to rank 5)

$A_1 = C_1$: original SW curve and massless monopole/dyon, F theory picture of D3/O7

Review of monodromy of $SU(2) = Sp(2)$ SW curve

How did they get massless monopole and dyon?

$$y^2 = (x^2 - u)^2 - \Lambda^4$$

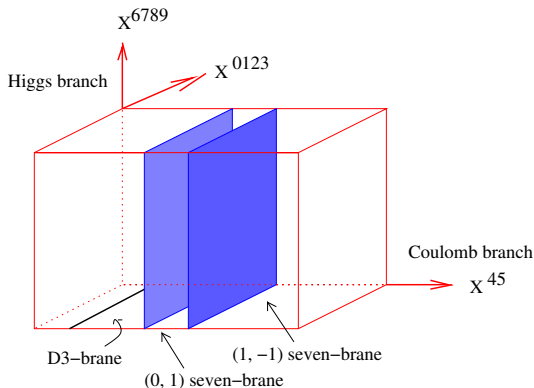


Start at a generic point in moduli space, then branch points on x -plane are separated. As you move around a singular locus in u -plane, a pair of branch points approach & change with each other. Their trajectory gives vanishing cycle.

$A_1 = C_1$: original SW curve and massless monopole/dyon, F theory picture of D3/O7

F-theoretic (quantum) 3/7 brane picture of $Sp(2)$

[Sen '96] [Banks-Douglas-Seiberg '96] & [Vafa 96]



Two (p, q) 7-branes are monopole and dyon in the original SW theory.

- Put r D3-branes as probes to get $Sp(2r)$ gauge theory.

[Douglas Lowe Schwarz 96]

SW curve for pure $SU(r+1) = A_r$

From [Klemm Lerche Yankielowicz Theisen 94]

$$\begin{aligned} y^2 = f_{SU(r+1)} &\equiv (x^{r+1} + u_1 x^{r-1} + u_2 x^{r-2} + \cdots + u_r)^2 - \Lambda^{2r+2} \\ &= f_+ f_- \end{aligned}$$

$$f_{\pm} \equiv x^{r+1} + u_1 x^{r-1} + u_2 x^{r-2} + \cdots + u_r \pm \Lambda^{r+1}$$

$$f_+ \equiv \prod_{i=0}^r (x - P_i) \quad f_- \equiv \prod_{i=0}^r (x - N_i)$$

Note that f_{\pm} do not share roots ($f_+ - f_- = 2\Lambda^{r+1} \neq 0$).

On the x -plane, only P_i 's (or N_i 's) can collide among themselves.

→ Discriminant factorizes $\Delta_x f_{SU(r+1)} = \# (\Delta_x f_+) (\Delta_x f_-)$

SW curve for pure $Sp(2r) = C_r$

By taking no-flavor limit of [Argyres Shapere 95], obtain

$$\begin{aligned} y^2 = f_{Sp(2r)} &\equiv \left(\prod_{a=1}^r (x - \phi_a^2) \right) \left(x \prod_{a=1}^r (x - \phi_a^2) + \Lambda^{2r+2} \right) \\ &= f_C f_Q \end{aligned}$$

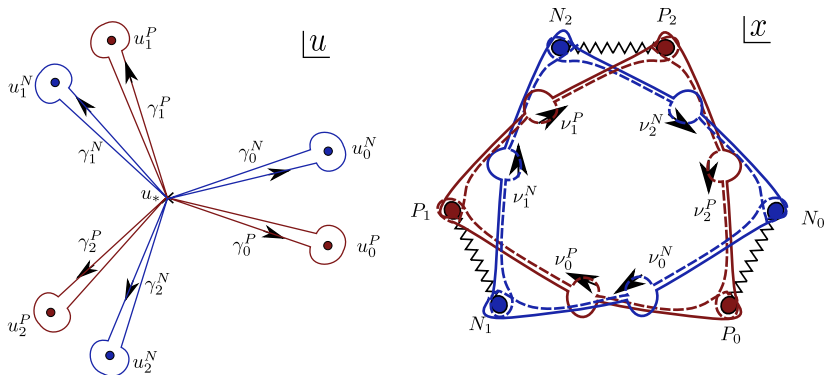
$$\begin{aligned} f_C &\equiv \prod_{a=1}^r (x - \phi_a^2) = \prod_{i=1}^r (x - C_i) \\ &= x^r + u_1 x^{r-1} + u_2 x^{r-2} + \cdots + u_r \\ f_Q &\equiv f_C x + \Lambda^{2r+2} = \prod_{i=0}^r (x - Q_i) \end{aligned}$$

Similarly as in $SU(r+1)$ curve, $\Delta_x f_{Sp(2r)} = \# (\Delta_x f_C) (\Delta_x f_Q)$

Identify massless dyons of A_r and C_r curves at $\Delta_x f = 0$

Vanishing cycles of SW curve for pure $SU(r+1)$

$SU(3)$ [Klemm Lerche Yankielowicz Theisen 94]



Identify massless dyons of A_r and C_r curves at $\Delta_x f = 0$

For $SU(r+1)$, in a region of moduli space given by

$$u_1 = \cdots = u_{r-2} = 0$$

$$u_r / \Lambda^{r+1} \in \mathbb{I},$$

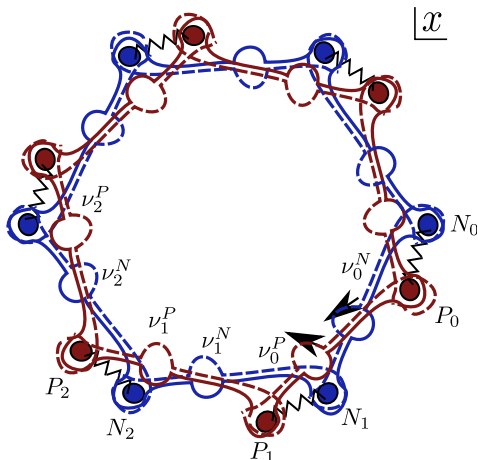
monodromy satisfies

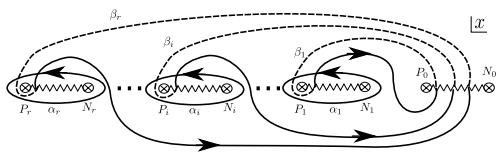
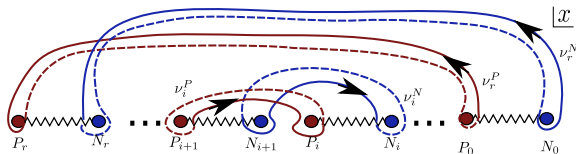
$$\nu_i^P \cap \nu_{i+1}^P = \nu_i^N \cap \nu_{i+1}^N = 1$$

$$\nu_i^P \cap \nu_i^N = -2$$

$$\nu_i^P \cap \nu_{i+1}^N = 2$$

all other intersection numbers vanish.



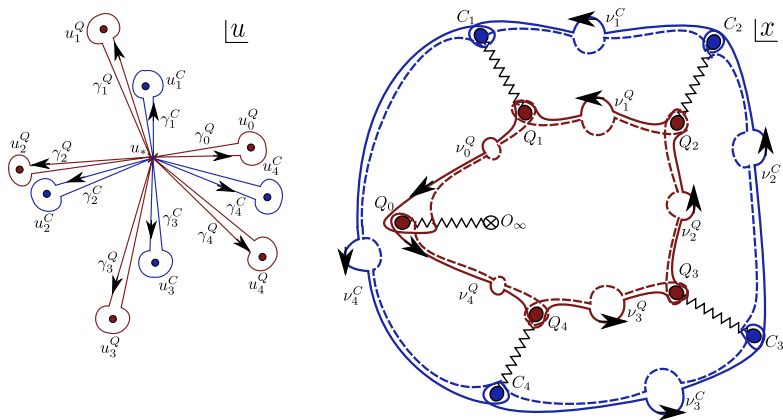
Identify massless dyons of A_r and C_r curves at $\Delta_x f = 0$ 

$$\begin{aligned} \nu_i^P &= \beta_{i+1} - \beta_i - \alpha_i \\ \nu_i^N &= \beta_{i+1} - \beta_i + \alpha_{i+1} - 2\alpha_i \\ i &= 1, \dots, r-1 \end{aligned}$$

$$\begin{aligned} \nu_0^P &= \beta_1 \\ \nu_r^P &= -\beta_r + \sum_{i=1}^{r-1} \alpha_i \\ \nu_0^N &= \beta_1 + \sum_{i=1}^r \alpha_i + \alpha_1 \\ \nu_r^N &= -\beta_r - 2\alpha_r \end{aligned}$$

Identify massless dyons of A_r and C_r curves at $\Delta_x f = 0$

$Sp(2r)$ monodromy



In a moduli region by $u_2 = \dots = u_{r-1} = 0$ and $u_r = \text{const}$: choose small enough u_r to keep $\nu_Q \cap \nu_C = 0$ (checked up to rank 5).

Identify massless dyons of A_r and C_r curves at $\Delta_x f = 0$

With the same choice of symplectic basis as $SU(r+1)$ before, the vanishing cycles are written as

$$\nu_0^Q = \beta_1$$

$$\nu_r^Q = -\beta_r - \sum_{i=1}^r \alpha_i - \alpha_r$$

$$\nu_r^C = \beta_1 - \beta_r - \sum_{i=2}^r \alpha_i$$

$$\nu_i^Q = \beta_{i+1} - \beta_i - \alpha_i$$

$$\nu_i^C = \beta_{i+1} - \beta_i + \alpha_{i+1}$$

$$i = 1, \dots, r-1$$

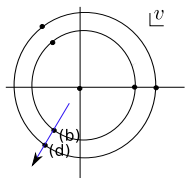
whose non-vanishing intersection numbers come from only

$$\nu_i^Q \cap \nu_{i+1}^Q = 1, \quad \nu_r^Q \cap \nu_0^Q = \nu_i^C \cap \nu_{i+1}^C = -1$$

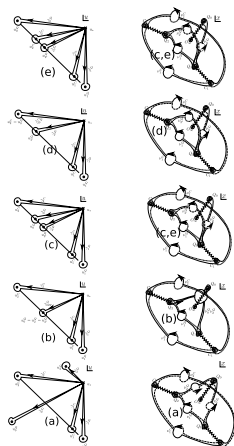
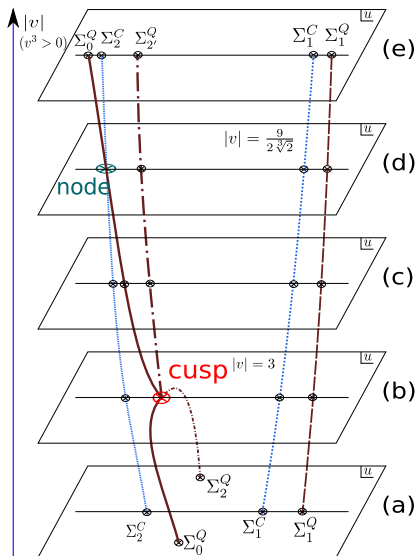
The spectra change as we move in moduli space. (C_2 example)

Example: Singularity structure of $Sp(4)=C_2$

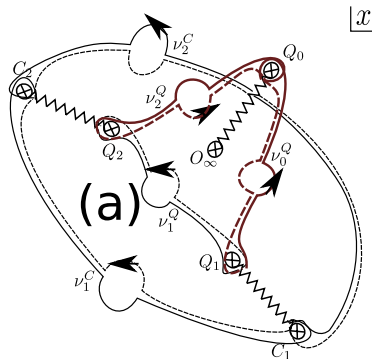
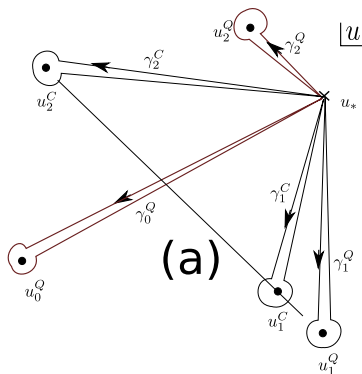
Example:
 $Sp(4) = C_2$



$\Delta_x f = 0$ at Σ
which
intersect at
 $\Delta_u \Delta_x f = 0$

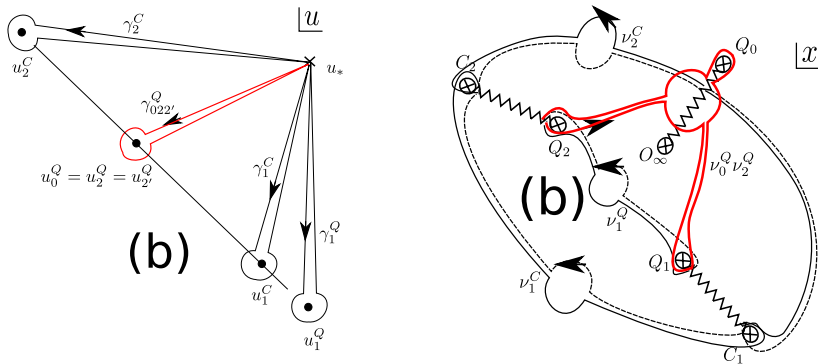


Example: Singularity structure of $Sp(4)=C_2$



Two cycles ν_0^Q and ν_2^Q vanish at two different moduli loci u_0^Q and u_2^Q respectively.

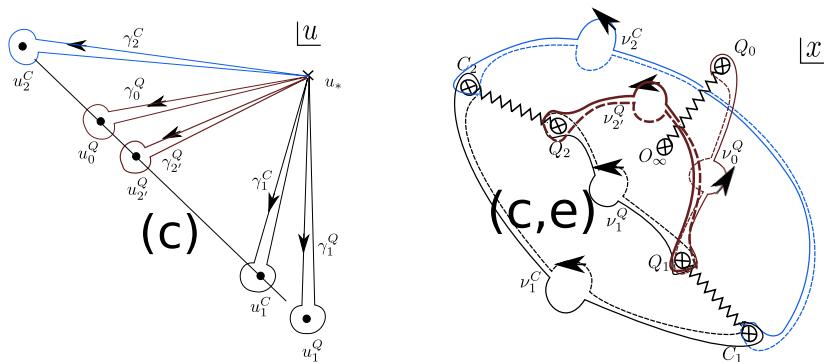
Example: Singularity structure of $Sp(4)=C_2$



As we vary in a moduli space (where u_0^Q and u_2^Q intersect tangentially), we hit a locus where they vanish simultaneously. The curve degenerates into $y^2 \sim (x - a)^3 \times \dots$. (Argyres-Douglas)

The red curve on the right does not give a well defined cycle.

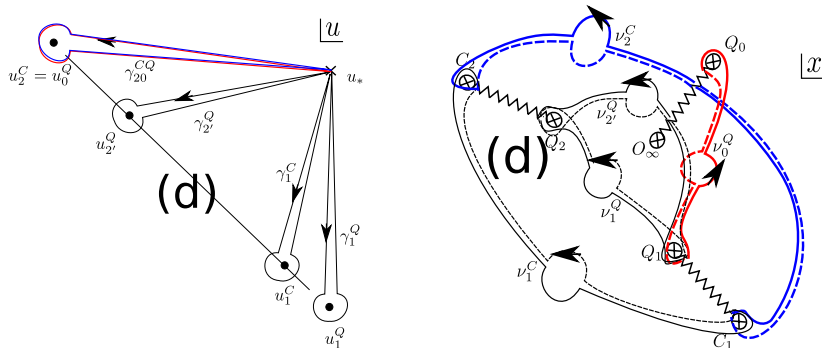
Example: Singularity structure of $Sp(4)=C_2$



Now u_0^Q and u_2^Q are separated, but u_0^Q runs toward another singularity locus u_2^C .

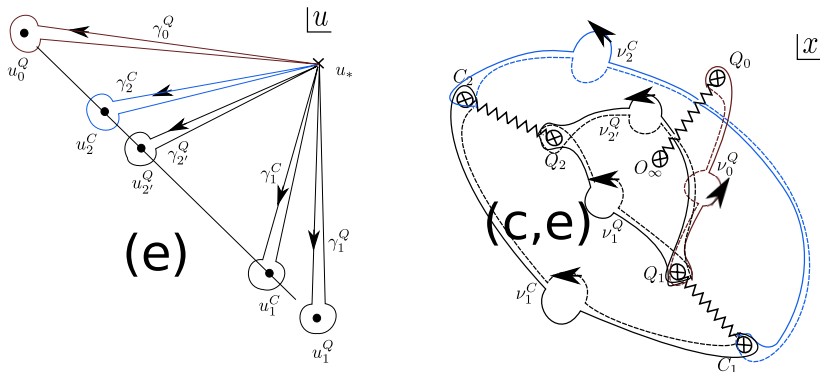
($\nu_2^Q = \nu_0^Q + \nu_2^Q$; Not related by symplectic transformation)

Example: Singularity structure of $Sp(4)=C_2$



This time, singularity loci u_0^Q and u_2^C intersect (node-like crossing). Vanishing cycles ν_0^Q and ν_2^C seem mutually non-local, but the curve looks like $y^2 = (x - a)^2(x - b)^2 \times \dots$ and it is not Argyres-Douglas form. (\therefore generalized AD)

Example: Singularity structure of $Sp(4)=C_2$



Now u_0^Q and u_2^C are separated, but the vanishing cycles did not change, unlike $(b \rightarrow c)$

$\Delta_u \Delta_x f$ captures higher singularity

Massless dyons coexist at a codim- $2_{\mathbb{C}}$ loci $\{\Delta_x f = \Delta_u \Delta_x f = 0\}$ [Argyres Plesser Seiberg Witten '95], where the curve looks like either of following two:

- ▶ $y^2 = (x - a)^3 \times \dots$
 - ▶ curve has **cusp**-like singularity (Argyres-Douglas)
 - ▶ $\Delta_x f = 0$ locus also intersects at $\Delta_u \Delta_x f = 0$ (o.o.v ≥ 3) with **cusp**-like singularity (tangential intersection)
 - ▶ two massless dyons are mutually non-local.
- ▶ $y^2 = (x - a)^2(x - b)^2 \times \dots$
 - ▶ curve has **node**-like singularity
 - ▶ $\Delta_x f = 0$ locus also intersects at $\Delta_u \Delta_x f = 0$ (o.o.v ≤ 2) with **node**-like singularity
 - ▶ two massless dyons are mutually local, with some exception (*generalized AD*: non-Argyres-Douglas & non-local loci) occurring at $SU(5 \uparrow), Sp(4 \uparrow)$

$\Delta_u \Delta_x f$ captures higher singularity

Generalized Argyres-Douglas loci

curve degen.n	$y^2 = (x - a)^3 \times \dots$	$y^2 = (x - a)^2(x - b)^2 \dots$	
o.o.v of $\Delta_u \Delta_x$	3	2	
shape of curve	cuspid	node	
shape of $\Delta_x = 0$	cuspid	node	
intersection	mutually non local	local	non-local
name	Argyres-Douglas	ML	gen AD

- ▶ generalized Argyres-Douglas: non-Argyres-Douglas & non-local loci (\exists conformal limit?)

Conclusion

- ▶ Discriminant $\Delta_x f = 0$ of the SW curve $y^2 = f$: identified all $2r + 1$ and $2(r + 1)$ massless dyons for $Sp(2r)$ and $SU(r + 1)$
- ▶ At double discriminant $\Delta_u \Delta_x f = 0$: massless dyons coexist. If order of vanishing ≥ 3 , then Argyres-Douglas. (checked up to rank 5)

Questions still remain...

- ▶ Behaviour in the Argyres-Douglas neighborhood?
- ▶ Rank r curve classification as rank 2 of [Argyres Crescimanno Shapere Wittig '05][Argyres Wittig '05]
- ▶ Global behaviour in moduli space: wall crossing [ShapereVafa99, GaiottoMooreNeitzke09]